

U.S. DEPARTMENT OF COMMERCE
National Technical Information Service

AD-A029 516

**Penetration of Electromagnetic
Pulses through Larger Apertures
in Shielded Enclosures**

Dikewood Industries, Inc.

May 1976

AFWL-TR-75-95

AFWL-TR-
75-95

ADA 029516



PENETRATION OF ELECTROMAGNETIC PULSES THROUGH LARGER APERTURES IN SHIELDED ENCLOSURES

University of Illinois
Urbana, IL 61801

for Dikewood Corporation
Albuquerque, NM 87106

May 1976

Final Report

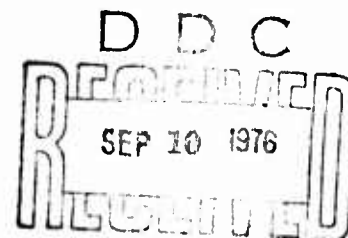
Approved for public release; distribution unlimited.

This research was sponsored by the Defense Nuclear Agency
under Subtask EB088, Work Unit 33, Work Unit Title:
Coupling Characteristics of Apertures.

Prepared for
Director
DEFENSE NUCLEAR AGENCY
Washington, DC 20305

AIR FORCE WEAPONS LABORATORY
Air Force Systems Command
Kirtland Air Force Base, NM 87117

REPRODUCED BY
NATIONAL TECHNICAL
INFORMATION SERVICE
U. S. DEPARTMENT OF COMMERCE
SPRINGFIELD, VA. 22161



This final report was prepared by the University of Illinois, Urbana, IL, for Dikewood Corporation, Albuquerque, NM, under Contract F29601-74-C-0010, Job Order WDNE 2705, with the Air Force Weapons Laboratory, Kirtland AFB, NM. Mr. William D. Prather (ELP) was the Laboratory Project Officer-in-Charge.

When US Government drawings, specifications, or other data are used for any purpose other than a definitely related Government procurement operation, the Government thereby incurs no responsibility nor any obligation whatsoever, and the fact that the Government may have formulated, furnished, or in any way supplied the said drawings, specifications, or other data is not to be regarded by implication or otherwise as in any manner licensing the holder or any other person or corporation or conveying any rights or permission to manufacture, use, or sell any patented invention that may in any way be related thereto.

This report has been reviewed by the Information Office (OI) and is releasable to the National Technical Information Service (NTIS). At NTIS, it will be available to the general public, including foreign nations.

This technical report has been reviewed and is approved for publication

William D. Prather
WILLIAM D. PRATHER
Project Officer

FOR THE COMMANDER

Larry W. Wood
LARRY W. WOOD
Lt Colonel, USAF
Chief, Phenomenology and
Technology Branch

James L. Griggs, Jr.
JAMES L. GRIGGS, JR.
Colonel, USAF
Chief, Electronics Division

EXEMPTION FOR	
White Section	Grey Section
UNCLASSIFIED	
JUSTIFICATION	
BY	
DISTRIBUTION/AVAILABILITY CODES	
Dist.	AVAIL. AND OF ST.
A	

DO NOT RETURN THIS COPY. RETAIN OR DESTROY.



UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER AFWL-TR-75-95	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) PENETRATION OF ELECTROMAGNETIC PULSES THROUGH LARGER APERTURES IN SHIELDED ENCLOSURES		5. TYPE OF REPORT & PERIOD COVERED Final Report
7. AUTHOR(s) R. Mittra L. Wilson Pearson University of Illinois		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS Dikewood Corporation Albuquerque, New Mexico 87106		8. CONTRACT OR GRANT NUMBER(s) F29601-74-C-0010
11. CONTROLLING OFFICE NAME AND ADDRESS Director Defense Nuclear Agency Washington, DC 20305		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS 62707H; WDNE2705
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office) Air Force Weapons Laboratory (ELP) Kirtland AFB, New Mexico 87117		12. REPORT DATE May 1976
		13. NUMBER OF PAGES 86 79
		15. SECURITY CLASS (of this report) UNCLASSIFIED
		15a. DECLASSIFICATION/DOWNGRADING SCHEDULE
16. DISTRIBUTION STATEMENT (of this Report) Approved for public release; distribution unlimited.		
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES This research was sponsored by the Defense Nuclear Agency under Subtask EB088, Work Unit 33, Work Unit Title: Coupling Characteristics of Apertures.		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Electromagnetic Fields and Waves Interaction and Coupling Aperture Penetration		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) The results of an initial investigation of the Singularity Expansion Method representation of the electromagnetic coupling through a rectangular aperture in a perfectly conducting sheet are reported. The problem is formulated in terms of the coupled Hallen-type integral equations for the dual problem of a rectangular plate. The integral equations are converted to a system of linear algebraic equations by way of the method of moments with subsectionally con- stant expansion functions and collocation testing. Several techniques used in		

DD FORM 1473 1 JAN 73 EDITION OF 1 NOV 65 IS OBSOLETE

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

PRICES SUBJECT TO CHANGE

UNCLASSIFIED

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

20. ABSTRACT (Continued)

minimizing the execution time of the computations are described. Some difficulties in accurately approximating the singularities of the system of integral equations by the singularities of the algebraic system are discussed. These difficulties arise because the subsectionally constant representation for the current cannot adequately represent the correct edge singularities in the currents on the plate. A set of pole trajectories indicative of the trends in pole location for the plate is reported. A listing of the pertinent computer code is provided.

UNCLASSIFIED

///

SECURITY CLASSIFICATION OF THIS PAGE(When Data Entered)

CONTENTS

<u>Section</u>		<u>Page</u>
I	INTRODUCTION	5
II	THIN-PLATE INTEGRAL EQUATION FORMULATION FOR COMPLEX WAVENUMBER	7
III	SYMMETRY CONDITIONS FOR THE NATURAL MODE CURRENTS	10
IV	THE NUMERICAL MODEL	16
V	ALGORITHMIC CONSIDERATIONS IN EVALUATING THE SYSTEM DETERMINANT	21
VI	NUMERICAL CHECKS ON THE ACCURACY OF THE POLES	25
VII	POLE TRAJECTORIES AS A FUNCTION OF SHAPE RATIO	31
VIII	CONCLUSIONS	33
	APPENDIX A: THE SELF-PATCH INTEGRATION	35
	APPENDIX B: THE SPARSE MATRIX ALGORITHMS	37
	APPENDIX C: PROGRAM LISTINGS	43
	REFERENCES	74

ILLUSTRATIONS

<u>Figure</u>		<u>Page</u>
1	Geometry of the Rectangular Plate	8
2	Lowest Order Natural Mode Current Pairs for Each of the Symmetry Cases, a) J_x Symmetric w.r.t. x-Axis and Symmetric w.r.t. y-Axis, b) Symmetric-Antisymmetric, c) Antisymmetric-Symmetric, and d) Antisymmetric-Antisymmetric	15
3	Subsectioning for the Discretization of the Integral Equations	17
4	Organization of the System of Linear Equations	19
5	a) Conceptual Zoning for Calculation of the Interaction Matrix, b) Example of the Four Interaction Contributions to a Single Source Term	22
6	Calculated Pole Locations for Thin-Strip for Varying Numbers of Zones in the x-Direction and Different Edge Treatments (Cylinder Results from Ref. 6)	26
7	Computer Pole Trajectory Under Varying w/L with Zoning Changes	28
8	Behavior of Singular Component of Current at the Edge Under Change in Transverse Zoning	30
9	Pole Trajectories as Computed with 4x3 zoning	32
B1	Submatrix Organization for the Sparse Matrix Algorithms, a) the Original Matrix, and b) Triangularized Form with the Generated CMAT4	38

TABLES

<u>Table</u>		<u>Page</u>
1	Compatible Current Symmetry Features	14
2	Matrix Formulation Parameters	20
B1	Primary Indexing Quantities in the Algorithm	42

SECTION I

INTRODUCTION

This report presents the results of an investigation for representing the transient electromagnetic coupling through a rectangular aperture by means of the singularity expansion method.

The singularity expansion method, which was introduced by Baum in 1971 (ref. 1), has been subsequently applied to many scatterer geometries. The essence of the singularity expansion method is the representing of the temporal response of a body in terms of the complex natural frequencies for the body.

Taylor et al. point out that the singularity expansion for an aperture in an infinite perfectly conducting screen can be determined in terms of that for the complementary perfectly conducting plate by way of Babinet's principle (ref. 2). This approach was taken in the work reported here. The remaining discussion is directed to the equivalent problem of determining the current distributions on the complementary plate geometry.

Rahmat-Samii and Mittra have derived a coupled pair of Hallén-type integral equations governing the current behavior on the rectangular plate (ref. 3). The work reported here builds on their work by generalizing the integral equations and solution method to the complex frequency plane for the

-
1. Baum, C. E., "On the Singularity Expansion Method for the Solution of Electromagnetic Interaction Problems," Interaction Note 88, Air Force Weapons Laboratory, Kirtland AFB, NM, December 1971.
 2. Taylor, C. D., Crow, T. T., and Chen, K-T, "On the Singularity Expansion Method Applied to Aperture Penetration: Part I Theory," Interaction Note 134, Air Force Weapons Laboratory, Kirtland AFB, NM, May 1973.
 3. Rahmat-Samii, Y. and Mittra, R., "Integral Equation Solution and RCS Computation of a Thin Rectangular Plate," Interaction Note 156, Air Force Weapons Laboratory, Kirtland AFB, NM, December 1973.

SEM application. The same method-of-moments formulation, as described in (ref. 3), is used, i.e., two-dimensional pulse expansion functions with collocation testing.

In order that the computation time be practical in a problem of this complexity, a great deal of care was given to algorithmic streamlining in the conduct of this work. The streamlining includes maximum exploitation of geometric symmetry, organization of calculations to make use of redundant terms and partial terms occurring in the calculation, and direct algorithmic exploitation of matrix sparseness. The end result is a highly efficient computer code. Key features of the algorithms are discussed in this report.

The pulse expansion appears to be inadequate in accurately modeling the rectangular plate. The difficulty, which relates to representing the actual size of the plate, is demonstrated and discussed herein. Remedies for the problem are suggested, but they are outside the scope of the present work.

By holding the zoning scheme invariant while the aspect ratio of the plate was changed, self-consistent pole trajectories for the four fundamental modes were determined. For the reasons cited above, these poles cannot claim to be the exact poles for the body. They are, however, indicative of the trends in the pole behavior for the plate under change in aspect ratio. These results are reported and discussed in this context.

SECTION II

THIN-PLATE INTEGRAL EQUATION FORMULATION FOR COMPLEX WAVENUMBER

Rahmat-Samii and Mittra (ref. 3) give an integral equation formulation for the rectangular plate subject to time-harmonic excitation. Their results may be directly extended to the complex wavenumber case. That is, for the geometry in Figure 1 with $\exp[st]$ time dependence, $s = \sigma + j\omega$ complex, and an incident plane-wave magnetic field component

$\vec{H}^i = [H_{ox}^i \hat{u}_x + H_{oy}^i \hat{u}_y + H_{oz}^i \hat{u}_z] \exp[j(k_x x + k_y y + k_z z)]$ the following coupled integral equations result:

$$\begin{aligned} \int_{-L/2}^{L/2} \int_{-w/2}^{w/2} \begin{Bmatrix} J_x(x,y) \\ J_y(x,y) \end{Bmatrix} K(x,y|x',y') dx' dy' &= \frac{1}{k_z} \begin{Bmatrix} H_{og}^i \\ -H_{ox}^i \end{Bmatrix} \exp[j(k_x x + k_y y)] \\ + \frac{\pi}{k} \begin{Bmatrix} -1 \\ -j \end{Bmatrix} \sum_{n=-\infty}^{\infty} C_n [j^{n+1} \exp[j(n+1)\phi] J_{n+1}(-s\rho/c) \\ + \begin{Bmatrix} 1 \\ -1 \end{Bmatrix} j^{n-1} \exp[j(n-1)\phi] J_{n-1}(-s\rho/c)] \end{aligned} \quad (1)$$

The kernel is given by

$$K(x,y | x',y') = \exp[-sR/c]/R \quad (2)$$

with

$$R = [(x - x')^2 + (y - y')^2]^{1/2}$$

The $J_x(x,y)$ and $J_y(x,y)$ denote the respective x and y components of current on the plate; $J_n(\zeta)$ denotes the Bessel function of the first kind; the C_n are unknown constants; c is the velocity of light; and (ρ, ϕ) are the polar coordinates for the point (x,y) on the plate. Equation (1) holds for $x \in (-L/2, L/2)$ and $y \in (-w/2, w/2)$, and $z = 0$.

It is pointed out that the two integral equations represented by (1) are

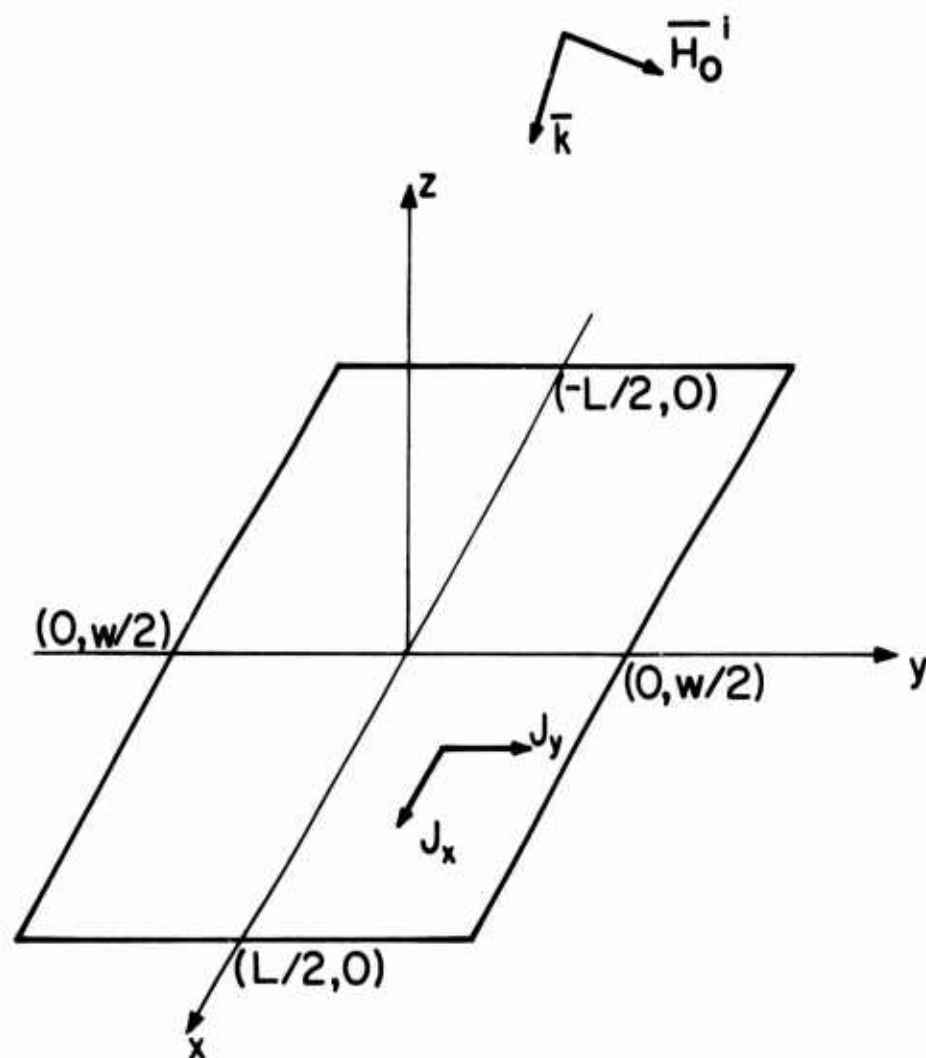


Figure 1. Geometry of the Rectangular Plate

coupled through the C_n in the summation in the right-hand side. This summation is simply a Bessel function expansion of the homogeneous solution to the wave equation which occurs in the derivation of (1). Details of arriving at this expansion are found in (ref. 3). The pair of integral equations (1) is complete in the sense that no additional constraints are needed to correctly specify the currents. It is noteworthy, however, that current solutions to (1) satisfy the Meixner's edge condition (ref. 4). Namely,

$$\left. \begin{aligned} J_x[\pm(L/2 - d), y] &\rightarrow d^{1/2} \\ J_y[\pm(L/2 - d), y] &\rightarrow d^{-1/2} \\ J_x[x, \pm(w/2 - d)] &\rightarrow d^{-1/2} \\ J_y[x, \pm(w/2 - d)] &\rightarrow d^{1/2} \end{aligned} \right\} \quad d \rightarrow 0 \quad (3)$$

The ability of a numerical solution to approximate the behavior of eqn. (3) is a key point in a subsequent discussion.

4. Bladel, J. Van, Electromagnetic Fields, McGraw-Hill, New York, pp. 385-387, 1964.

SECTION III

SYMMETRY CONDITIONS FOR THE NATURAL MODE CURRENTS

The natural frequencies of (1) occur when the complex frequency s is such that there are non-trivial J_x and J_y and the accompanying C_n satisfying (1) for $\bar{H}^1 = 0$. Such J_x and J_y solutions are natural mode current solutions for the rectangular plate, and the concomitant value of s is a pole of the plate. The vanishing of incident wave dependence gives rise to symmetry in the integral equations. By discerning the symmetry relations a priori and bringing them to bear upon solution procedures, one gains significant computational savings in the numerical solution for poles and natural modes. These symmetry relations are explored in this section.

The excitation-free form of (1) is

$$\int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_x K(x,y|x',y') dx' dy' = \frac{j\pi c}{s} \sum_{-\infty}^{\infty} C_n \left\{ j^{n+1} \exp[j(n+1)\phi] J_{n+1}(-sp/c) \right. \\ \left. + j^{n-1} \exp[j(n-1)\phi] J_{n-1}(-sp/c) \right\} \quad (4a)$$

and

$$\int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_y K(x,y|x',y') dx' dy' = \frac{\pi c}{s} \sum_{-\infty}^{\infty} C_n \left\{ j^{n+1} \exp[j(n+1)\phi] J_{n+1}(-sp/c) \right. \\ \left. - j^{n-1} \exp[j(n-1)\phi] J_{n-1}(-sp/c) \right\} \quad (4b)$$

By using the symmetry of the Bessel function with respect to order, expanding the exponentials by way of Euler's identity, and appropriately adjusting the indices, one arrives at the following equation after some manipulation.

$$\begin{aligned}
& \int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_x K \, dx' \, dy' \\
&= \frac{j\pi c}{s} \sum_{n=0}^{\infty} \left\{ j^{n+1} d_n^+ [\cos(n+1)\phi J_{n-1}(-s\rho/c) - u_{n-1} \cos(n-1)\phi J_{n-1}(-s\rho/c)] \right. \\
&\quad \left. - j^n d_n^- [\sin(n+1)\phi J_{n+1}(-s\rho/c) - \sin(n-1)\phi J_{n-1}(-s\rho/c)] \right\} \quad (5a)
\end{aligned}$$

and

$$\begin{aligned}
& \int_{-L/2}^{L/2} \int_{-w/2}^{w/2} J_y K \, dx' \, dy' \\
&= \frac{j\pi c}{s} \sum_{n=0}^{\infty} \left\{ j^{n+1} d_n^+ [\sin(n+1)\phi J_{n+1}(-s\rho/c) + \sin(n-1)\phi J_{n-1}(-s\rho/c)] \right. \\
&\quad \left. + j^n d_n^- [\cos(n+1)\phi J_{n+1}(-s\rho/c) + u_{n-1} \cos(n-1)\phi J_{n-1}(-s\rho/c)] \right\} \quad (5b)
\end{aligned}$$

where

$$d_n^{\pm} = C_n \pm C_{-n}$$

and

$$u_n = \begin{cases} 1, & n \geq 0 \\ 0, & n < 0 \end{cases}$$

It is noted that the d_n^+ multiply terms containing cosine functions in the J_x equation, while they multiply terms containing sine functions in the J_y equation. The situation is reversed for the d_n^- .

Because of the symmetry properties of the kernel, the integral operator on the left-hand sides of (5) produces a function whose symmetry character is identical to that of the current on which it operates. Then, for a given current symmetry, only part of the d_n^{\pm} on the right-hand side may be non-zero because of the symmetries possessed by the trigonometric terms. Thus, the respective symmetries for J_x and J_y , which are compatible, and the

surviving terms in the right-side series may be discerned by 1) postulating a symmetry for J_x , 2) determining from (5a) which right-hand side terms survive so as to be compatible with the J_x symmetry, 3) observing in (5b) the variation which terms have non-zero coefficients, and 4) determining the J_y symmetry conditions compatible with the postulated J_x symmetry conditions.

For example, if J_x is symmetric with respect to the y axis and anti-symmetric with respect to the x axis, only $\sin(n+1)\phi$ terms with n even are compatible in (5a). Thus, only d_n^- , n even, may be non-zero. In the right-hand side of (5b), the coefficients multiply $\cos(n+1)\phi$ terms with n even. These cosines sum to functions which are antisymmetric with respect to the y axis and symmetric with respect to the x axis. Stated mathematically, if

$$J_x(x,y) = J_x(-x,y) \quad (6a)$$

and

$$J_x(x,y) = -J_x(x,-y) \quad (6b)$$

then

$$d_n^+ = 0, \quad \text{for all } n, \quad (6c)$$

$$d_n^- = 0, \quad n \text{ odd}, \quad (6d)$$

and

$$J_y(x,y) = -J_y(-x,y) \quad (6e)$$

$$J_y(x,y) = J_y(x,-y) \quad (6f)$$

These vector symmetries are in accord with the general symmetry relations given by Baum (ref. 5). The information in (6) may be used to reduce the complexity of the integral equations (4), viz., by (6a,b,e,f) the range of each integration may be halved while by (6c,d) the zero terms of the right-hand side are known a priori:

$$\begin{aligned} & \int_0^{L/2} \int_0^{w/2} J_x K^+(x,y|x',y') dx' dy' \\ &= \frac{\pi c}{s} \sum_{\substack{n=0 \\ n \text{ even}}}^{\infty} d_n^- j^{n-1} [\sin(n+1)\phi J_{n-1}(-sp/c) - \sin(n-1)\phi J_{n-1}(-sp/c)] \end{aligned} \quad (7a)$$

-
5. Baum, C. E., "Interaction of Electromagnetic Fields with any Object which has an Electromagnetic Symmetry Plane," Interaction Note 63, Air Force Weapons Laboratory, Kirtland AFB, NM, March 1971.

and

$$\int_0^{L/2} \int_0^{w/2} J_y K^+(x,y|x',y') dx' dy' \\ = \frac{\pi c}{s} \sum_{\substack{n=0 \\ n \text{ even}}}^{\infty} j^{n+1} d_x^- [\cos(n+1)\phi J_{n+1}(-s\rho/c) + u_{n-1} \cos(n-1)\phi J_{n-1}(-s\rho/c)] \quad (7b)$$

where

$$K^+(x,y|x',y') = K(x,y|x',y') - K(x,y|-x',y') \\ + K(x,y|x',-y') - K(x,y|-x',-y') \quad (8a)$$

and

$$K^{++}(x,y|x',y') = K(x,y|x',y') + K(x,y|-x',y') \\ - K(x,y|x',-y') - K(x,y|-x',-y') \quad (8b)$$

For subsequent reference

$$K^{++}(x,y|x',y') = K(x,y|x',y') + K(x,y|-x',y') \\ + K(x,y|x',-y') + K(x,y|-x',-y') \quad (8c)$$

and

$$K^{--}(x,y|x',y') = K(x,y|x',y') - K(x,y|-x',y') \\ - K(x,y|x',-y') + K(x,y|-x',-y') \quad (8d)$$

are defined as well. Equations (7) are enforced for $z = 0$, $x \in (0, L/2)$ and $y \in (0, w/2)$.

Table 1 summarizes the four symmetry cases which are derived as in the foregoing discussion. By means of this table, four integral equation pairs can be constructed in the spirit of (7) by replacing the kernels in (7) with the appropriate kernels from the table and retaining only the non-vanishing terms in the series expansion.

Figure 2 depicts qualitatively the respective modal current distributions for the lowest frequency natural resonance exhibiting each symmetry.

Table 1

COMPATIBLE CURRENT SYMMETRY FEATURES

J_x				J_y				
Sym. w.r.t. x axis	Sym. w.r.t. y axis	Kernel	Compatible Trig. Fns.	Coefs. $\neq 0$	Compatible Trig. Fns.	Kernel	Sym. w.r.t. x axis	Sym. w.r.t. y axis
sym	sym	K^{++}	$\cos 2n\phi$	d_{2n+1}^+	$\sin 2n\phi$	K^{--}	anti	anti
sym	anti	K^{+-}	$\cos (2n + 1)\phi$	d_{2n}^+	$\sin (2n + 1)\phi$	K^{-+}	anti	sym
anti	sym	K^{-+}	$\sin (2n + 1)\phi$	d_{2n}^-	$\cos (2n + 1)\phi$	K^{+-}	sym	anti
anti	anti	K^{--}	$\sin 2n\phi$	d_{2n+1}^-	$\cos 2n\phi$	K^{++}	sym	sym

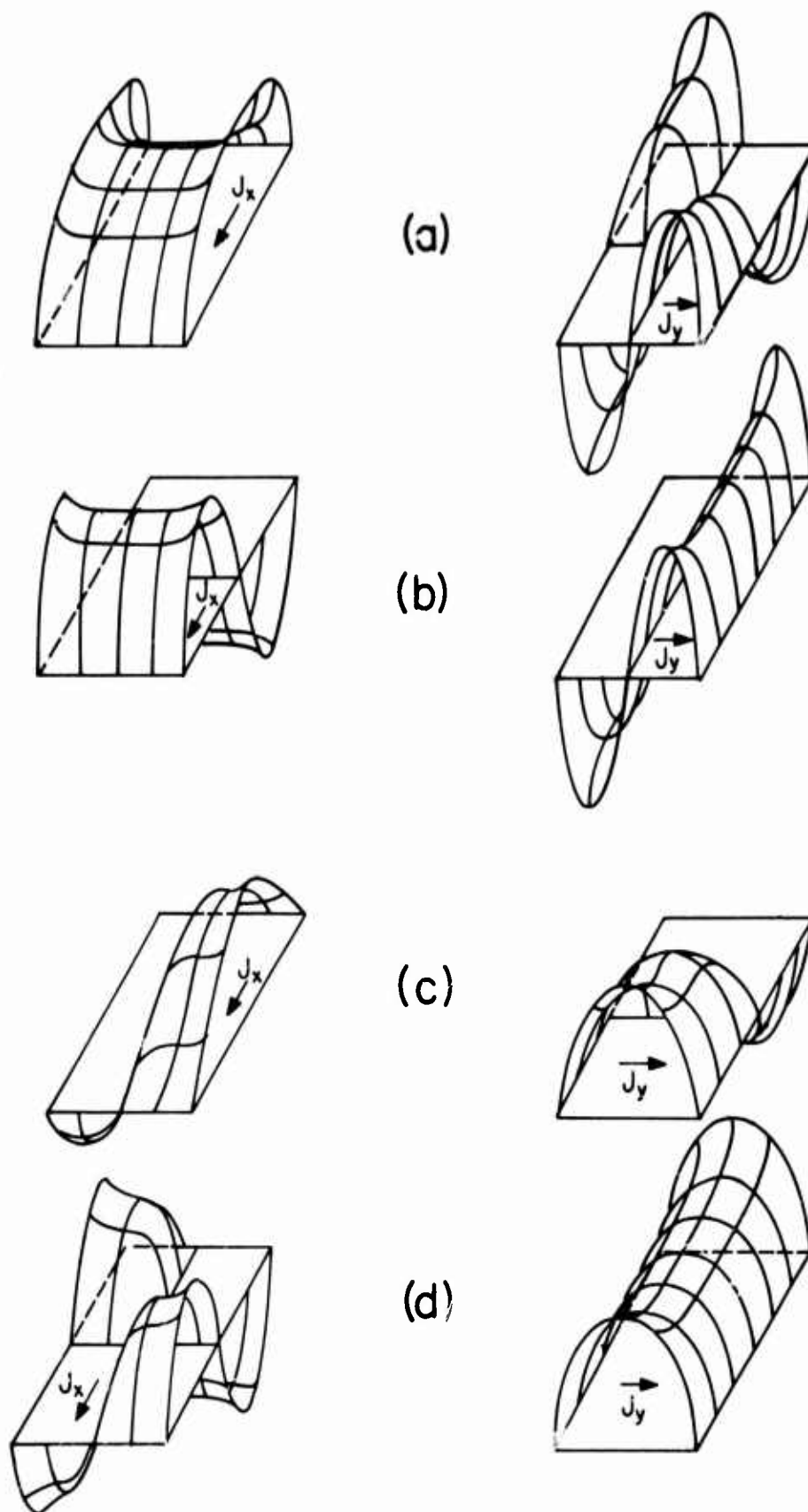


Figure 2. Lowest Order Natural Mode Current Pairs for Each of the Symmetry Cases, a) J_x Symmetric w.r.t. x-Axis and Symmetric w.r.t. y-Axis, b) Symmetric-Antisymmetric, c) Antisymmetric-Symmetric, and d) Antisymmetric-Antisymmetric

SECTION IV

THE NUMERICAL MODEL

The integral equation pair of the form (7) for each of the four symmetry cases can be discretized by the method of moments. In the work reported here, two-dimensional, subsectionally constant expansion functions were used with collocation testing. The zoning scheme is represented in Figure 3.

The unknown currents J_x and J_y were expanded in piecewise constant functions as in (ref. 3) with subsectioning of the form given in Figure 3. Notice that half-width patches are used at the edges of the plate so that match points lie precisely on the edge of the plate. The half-width pulse has proved useful in realizing the actual electrical size of a body in one-dimensional problems (ref. 6). Some numerical experimentation was also done with full-sized pulses on the edges and comparative results are reported in a later section. Some difficulties occur in definition of the edge of the plate in the present formulation because of the presence of two current components which have the asymptotic behavior given in (3). This difficulty is discussed in a later section.

The boundary condition $J_{\text{norm}} = 0$ must be enforced on selected patches at the edge of the plate as discussed in (ref. 3). Concomitantly, only as many d_n^{\pm} 's are retained in the right-hand side summation in (7) as there are current values preassigned to zero. The shaded patches in Figure 3 indicate the selection of patches where a current component is preassigned a zero value. At the corner patch, both components are preassigned zero values.

-
6. Butler, C. M., "Integral Equation Solution Methods," in "Wire Antennas and Scatterers," Short Course Notes, University of Mississippi, April 1972. (See also IEEE Trans., v. AP-20, pp. 731-736, 1972.)

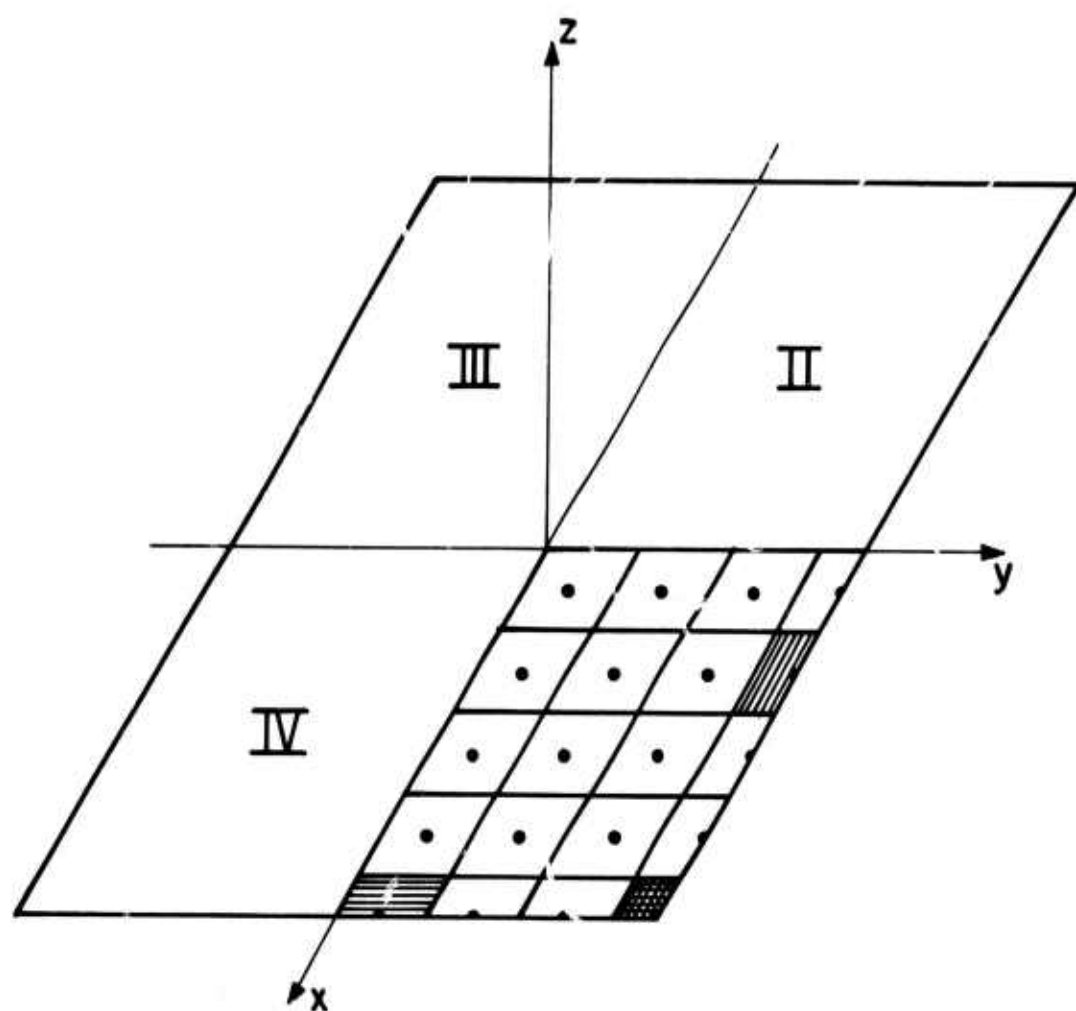


Figure 3. Subsectioning for the Discretization of the Integral Equations

By assigning one match point per expansion patch and by retaining one series expansion term for each current value preassigned in each of the two integral equations, a consistent (i.e. square) system of linear equations results. The truncated summation is taken to the left-hand side so that a homogeneous system results. The matrix organization used to represent these equations is given in Figure 4. Table 2 defines the computer variables noted in Figure 4, primarily for reference purposes in the next section.

The matrix that results is a function of the complex frequency s . A natural resonance occurs when s has a value such that the matrix has a zero determinant; hence, the homogeneous system of equations has a non-trivial solution. The next section explores some algorithmic considerations in the efficient generation and manipulation of the matrix.

$$\begin{bmatrix} \begin{bmatrix} M_x \\ (NI1 \times NJ1) \end{bmatrix} \\ \begin{bmatrix} 0 \\ (NI2 \times NJ1) \end{bmatrix} \end{bmatrix} \begin{bmatrix} 0 \\ (NI1 \times NJ2) \\ M_y \\ (NI2 \times NJ2) \end{bmatrix} \begin{bmatrix} M_\Sigma \\ (NI1 + NI2 \times NPRES) \end{bmatrix} \begin{bmatrix} J_x \\ (NJ1) \\ J_y \\ (NJ2) \\ d \\ (NPRES) \end{bmatrix} = \begin{bmatrix} 0 \end{bmatrix}$$

Figure 4. Organization of the System of Linear Equations

Table 2

MATRIX FORMULATION PARAMETERS

NI1	Number of match points on the zoned quadrant of the plate.
NI2 = NI1	
NPRED	Number of patches along the $ x = L/2$ edge where J_x is preassigned to zero.
NPRED	Number of patches along the $ y = w/2$ edge where J_y is preassigned to zero.
NJ1 = NI1-NPRED	Number of unknown current values in J_x expansion.
NJ2 = NI2-NPRED	Number of unknown current values in J_y expansion.
$\left. \begin{array}{l} NJ3 \\ NPRED \end{array} \right\} = NPRED - NPRED$	Number of preassigned current values (Also the number of coefficients retained in summation).

SECTION V

ALGORITHMIC CONSIDERATIONS IN EVALUATING THE SYSTEM DETERMINANT

Some considerations taken into account in generating the system matrix and evaluating its determinant efficiently are discussed in this section. Since these two operations must be repeatedly carried out for many values of s in the course of determining the natural frequencies of the plate, it is essential that clean, efficient computer programming and coding be used so that execution of the program will be affordable. The volume of code in the algorithms is consistently compromised toward a larger size in order to meet the following two time-efficient objectives:

1. Avoidance of calculating the same quantity twice; and
2. Avoidance of logical decisions, particularly those which might be imbedded in loops.

The program is discussed in the context of the following major segments:

1. Computation of an "interaction matrix";
2. Construction of the non-zero submatrices of the system matrix from the interaction matrix;
3. Computation of the series terms' submatrix; and
4. Determinant evaluation.

The major contribution to the elimination of redundant calculations is the one-time computation of an "interaction matrix" which is made up of the individual kernel integral terms from (2) for all argument combinations which occur in the computation. The subsequent program step then picks, by subscript, entries from this matrix and constructs the appropriate kernel from one of equations (8) according to the symmetry conditions being solved. This procedure can be viewed in terms of the layout given in Figure 5a. The terms in the interaction matrix are those evaluated for the match-point as

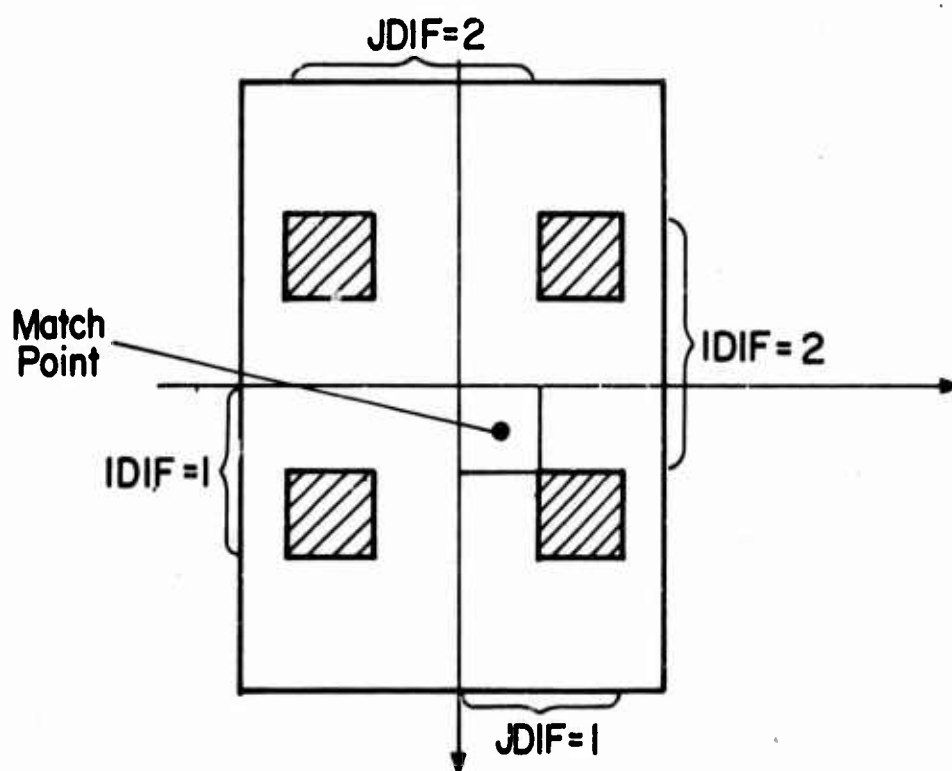
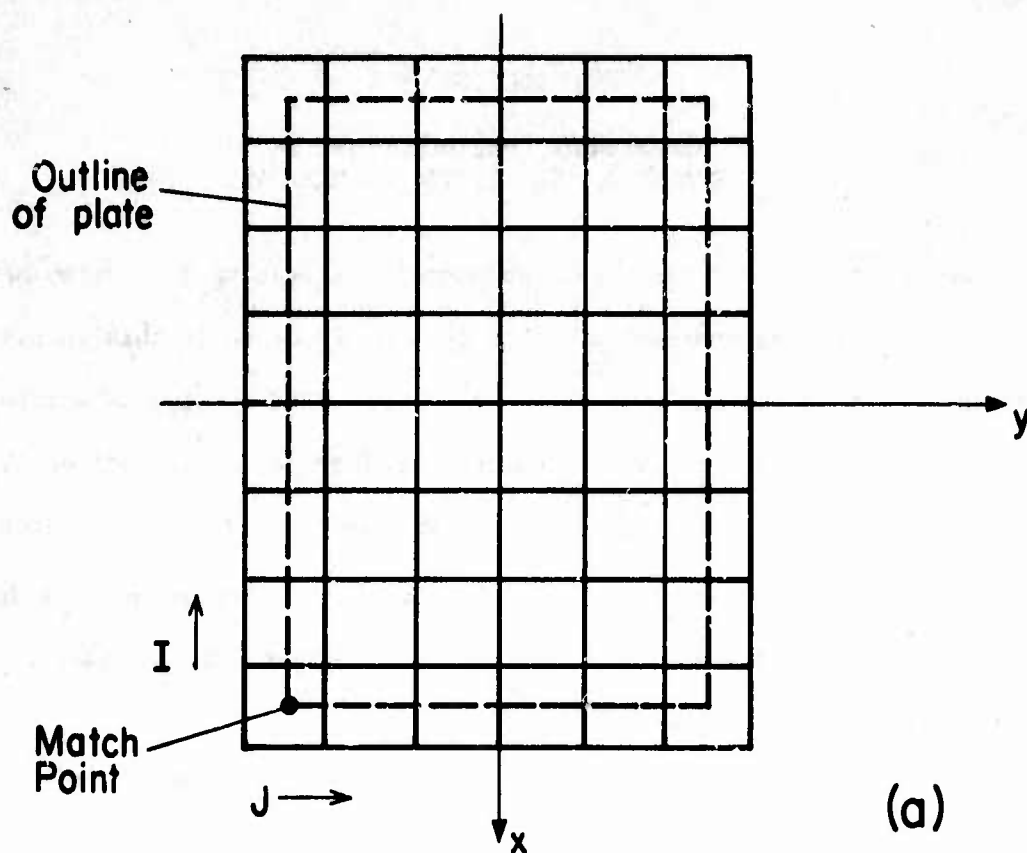


Figure 5. a) Conceptual Zoning for Calculation of the Interaction Matrix, b) Example of the Four Interaction Contributions to a Single Source Term

shown in the lower left with the source patches indexed over the entire plate to generate the matrix. Thus, all geometric relationships which occur in the kernel terms are encompassed in the calculation. Note that all source patches are full patches for this calculation. The effect of half patches at the edges is accounted for by weighting by a factor of $1/2$ the edge contributions. The kernel integral appropriate to the symmetry is constructed by summing with correct signs the appropriate elements from the matrix. Figure 5b gives an example of the four source patches entering into one kernel integral.

Differing degrees of sophistication are required in the calculation of the interaction terms depending on the spacing of the patches for which an interaction is being calculated. For the self patch, i.e., the patch in which the match-point resides, the integration of the kernel must be performed analytically because of the integrable singularity in the kernel there. Appendix A gives a series approximation to this integral. The result in Appendix A is evaluated directly in the program. For the patches adjacent to the patch containing the match point, the kernel is a rapidly varying but well-behaved function. The integration over these patches is evaluated numerically by a polynomial approximation. For patches further separated, the kernel is slowly varying and the integral is evaluated approximately as the product of the value of the kernel at the center of the patch and the area of the patch.

Some minor time economy is achieved in the matrix filling algorithm, which is a four-dimensional loop. The economy is found in the form of decision-free indexing, that is, the source contributions from interior patches, from $|x| = L/2$ edge patches, from $|y| = w/2$ edge patches, and from corners take on different forms. Rather than index over all source patches

with logical decisions implemented to discriminate among the four cases above, four different loops are used.

The computation of the series term submatrix is relatively straightforward. Because the Bessel-trigonometric products appear in two terms each, they are all precalculated and stored in a vector. The individual terms are then constructed from them.

The determinant evaluation profits significantly from an exploitation of the sparceness of the matrix. Either of two approaches may be taken to the sparse matrix manipulations. One is to separate the matrix algebraically and calculate an inverse as a composite of inverses of terms involving the submatrices. The alternative approach is to attack the matrix directly with Gaussian elimination, and to exploit the sparceness directly in the algorithm. The latter approach was chosen for the present purpose because it is judged to be slightly faster computationally and because in order to determine natural mode solutions for the SEM formulation, the homogeneous system of equations occurring at a pole must be backsolved. The algorithm resulting from the second approach is described in Appendix B.

The determinant evaluation routine itself appears in Appendix C as the function routine CPLATE.

SECTION VI

NUMERICAL CHECKS ON THE ACCURACY OF THE POLES

The results of some numerical checks on the accuracy of the pole location as determined from the numerical model described in Sections II through V are reported. It is shown that the model predicts well the poles for narrow strips possessing essentially thin scatterer resonance properties. Difficulties occur, however, in obtaining self-consistent results under different zone sizes for plates with larger aspect ratios. It is conjectured that the difficulty occurs because the subsectionally constant current representation is inadequate to represent the correct edge behavior for the currents—particularly the singular behavior for the parallel component. The rationale behind this conjecture is discussed.

Initial tests on the accuracy of the model were made for a rectangular strip with a shape ratio $w/L = 1/10$. Such a strip has an approximate equivalent dipole whose diameter-to-length ratio is $1/10\pi$.

Figure 6 gives the results of pole determinations for the first two poles for various numbers of pulses in the expansion of the current and for two different treatments of the edge pulse. The strip was zoned with one pulse across a quadrant. The numbers indicated in the figure are the numbers of pulses along the longitudinal direction of a quadrant. The "half-pulse" results are those obtained by the zone scheme described in Section IV where half-width pulses are placed at the edge so that match points fall at the edge. The "full-pulse" results are those obtained by zoning the plate with equal-sized pulses over the entire quadrant. In the latter case an a posteriori adjustment is made in the data correcting the length of the strip such that the end of the corrected strip lies at the end match-point.

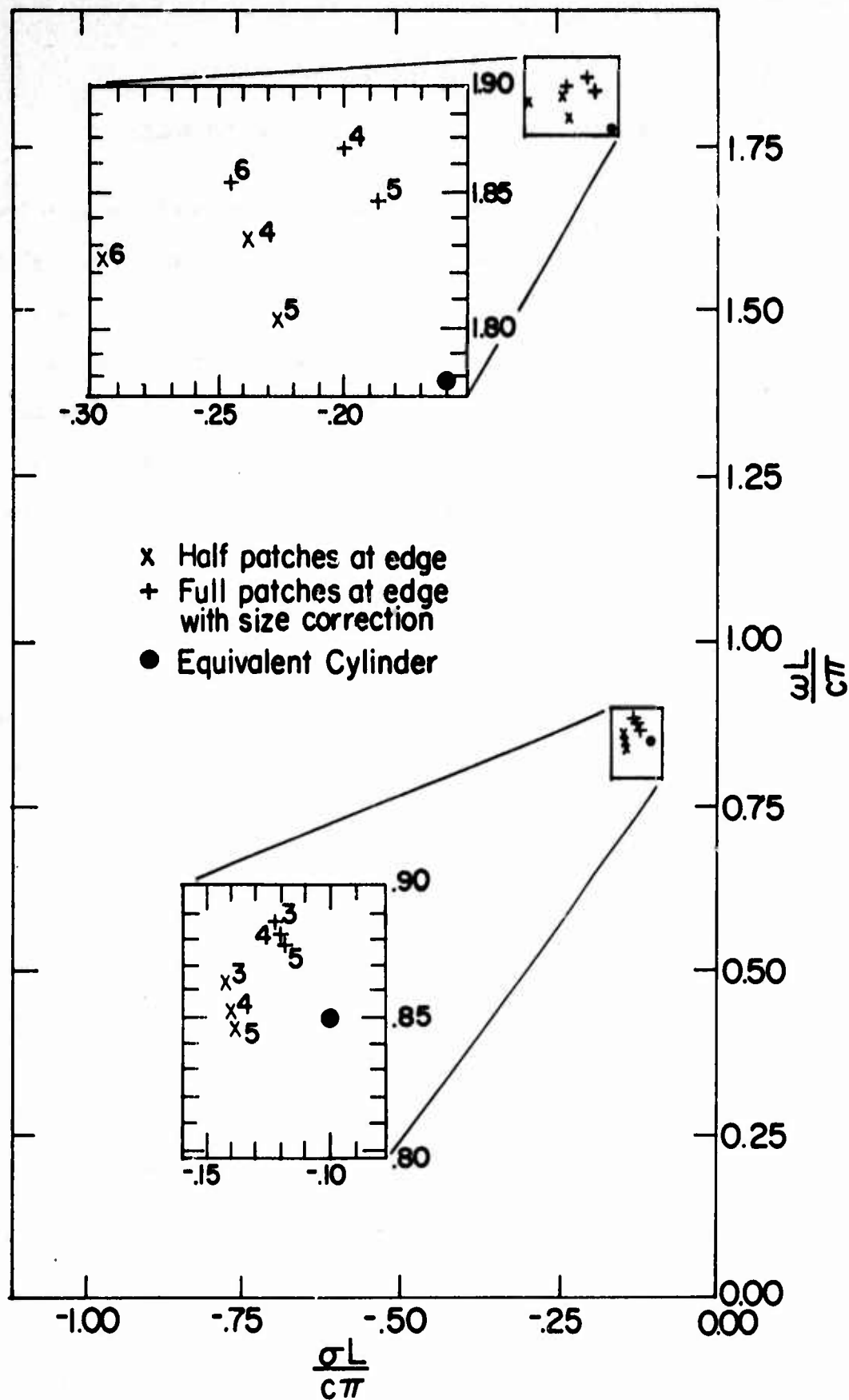


Figure 6. Calculated Pole Locations for Thin-Strip for Varying Numbers of Zones in the x-Direction and Different Edge Treatments (Cylinder Results from Ref. 6)

It is seen that the differences are small both for varying order and increasing pulse density. The $NX = 6$ results for the second pole show some departure from the trend established by the results for $NX = 4$ and $NX = 5$. This is attributable to the fact that the matrix is on the brink of numerical instability for $NX = 6$. The results for $NX = 7$, which are not shown, are observed to be meaningless because of the instability manifested.

For comparison purposes, the first two poles for an equivalent cylinder (one whose circumference equals the strip width) are given as found in ref. 7. These results are judged reliable inasmuch as they have been cross-checked among several integral equation formulations and for several method-of-moments schemes. The equivalent radius taken is, of course, an approximation. It is seen that the half-pulse solutions compare slightly more favorably with the cylinder results. Because of this, and moreover, because the a posteriori data adjustment is avoided with the half-pulse scheme, it was used consistently in the remaining solutions.

The stable convergence properties of the numerical model exhibited for the thin-strip solution are not manifested for higher aspect ratios. The reason for the difference is that the strip is essentially a one-dimensional problem and the transverse (y-directed) component of current has little effect on the dominant longitudinal current. For wider structures the coupling is significant and inadequacies in the modeling of the singularities of the currents produce inaccuracies which are highly sensitive to zoning.

Figure 7 shows the results obtained for a pole trajectory as a function of the shape factor w/L where the zoning in the y-direction was adjusted

-
7. Umashankar, K. R., "The Calculation of Electromagnetic Transient Currents on Thin Perfectly Conducting Bodies Using the Singularity Expansion Method," Ph.D. Thesis, University of Mississippi, pp. 33-34, August 1974, (See also Tesche, F. M., IEEE Trans., Vol. AP-21, No. 1, pp. 53-62, 1972.)

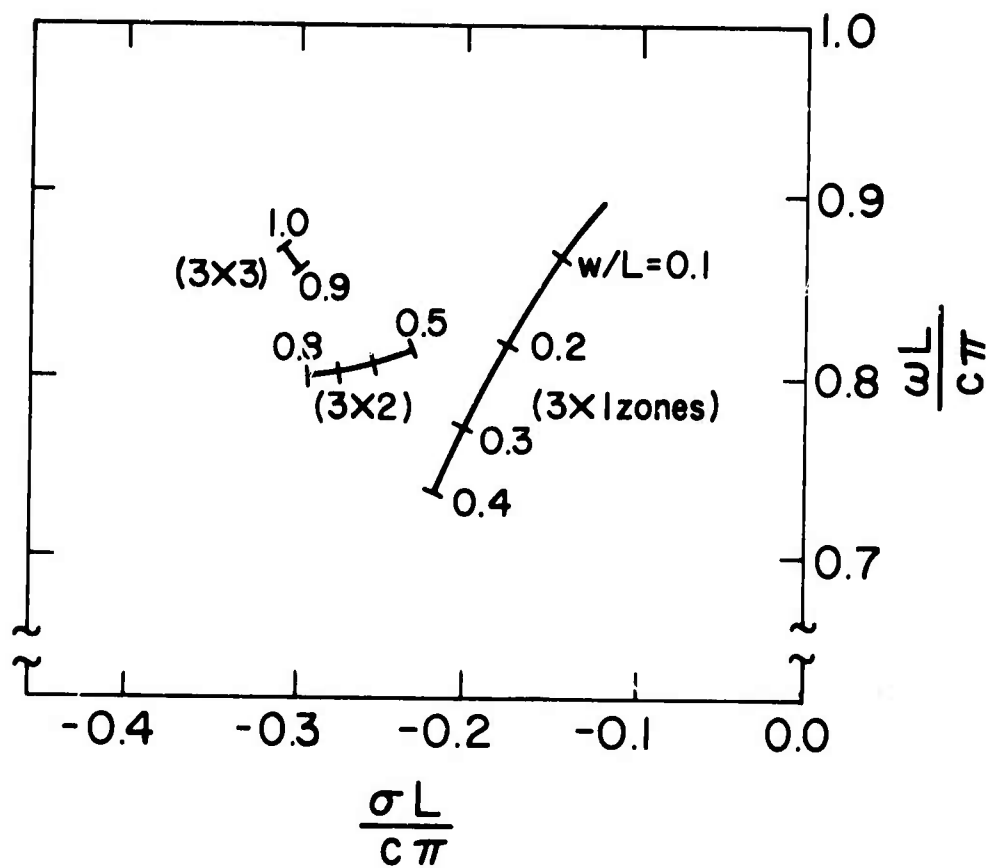


Figure 7. Computer Pole Trajectory Under Varying w/L with Zoning Changes

according to the value of w . It is evident that the solutions are unstable with respect to the zoning on the plate. Attempts to increase the number of zones significantly to improve upon the situation resulted in numerical instabilities in the matrix which cause the pole search iteration to fail.

The reason for the difficulty manifested in Figure 6 is believed to lie in the way that the edge of the plate is defined with the piecewise constant current expansion. Consider the characteristics of the two current components along a line traversing the plate in the y -direction as shown in Figure 8. The correct edge behavior at $|y| = w/2$ is that given in equations (3). The zoning scheme, however, forces $J_x(x, \pm w/2)$ to take a finite value. The current extrapolates to a singular point for some $y > w/2$, i.e., the numerical model represents a plate whose width is greater than w .

If the number of transverse zones is increased as indicated by the dashed curve in Figure 8, the steepness of the edge behavior of J_x is increased, and the extrapolation is characteristic of a narrower plate as compared to the first case. This narrowing of the effective width in the model is reflected in an increased Q (resonance quality factor) as the jumps in Figure 7 indicate.

One is tempted to conclude that a full-width pulse at the edge is a cure for the problem since the point at which the pulse current is defined is shifted relative to the edge as zoning is changed with full-width pulses. The effect of this procedure is to transfer the problem from component of current whose behavior is singular at the edge to the component which goes to zero. With full pulses at the edges, the normal component of current would go to zero interior to the plate rather than at the edge as it properly should.

An approach which is potentially a remedy for this difficulty is discussed in the conclusions.

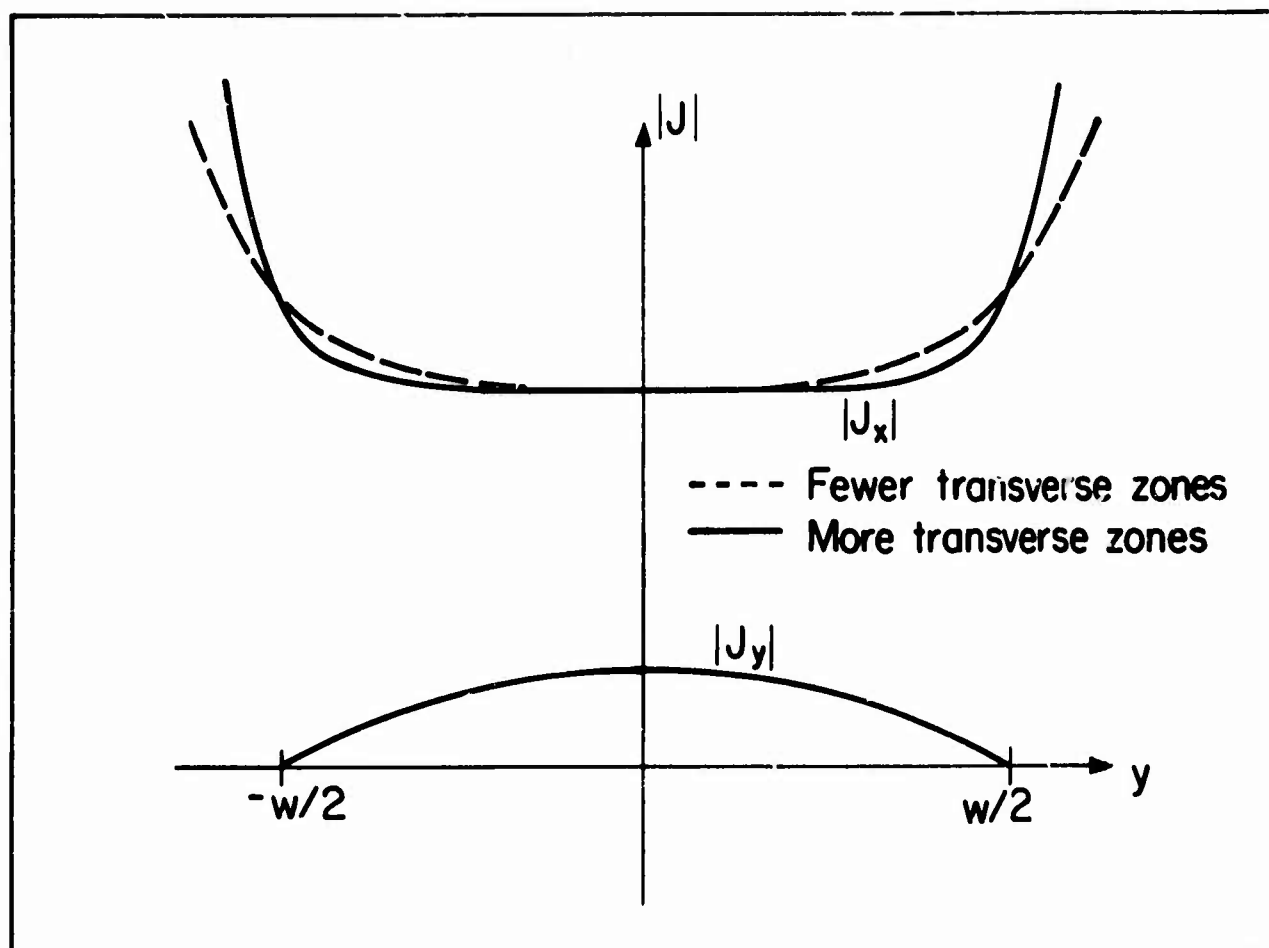


Figure 8. Behavior of Singular Component of Current at the Edge Under Change in Transverse Zoning

SECTION VII

POLE TRAJECTORIES AS A FUNCTION OF SHAPE RATIO

Figure 9 gives the results obtained for pole location for the lowest order pole of each of the symmetries as a function of w/L . Clearly, as the previous section indicates, the results cannot be taken as the correct results for the plate. However, the zoning was fixed for all w/L and the results are expected to reflect the proper trends for these trajectories.

The ++ and +- modes are in essence dipole modes, and their poles show the general lowering of Q as w/L increases. (The ++ indicates that the J_x component is symmetric both with respect to the x and y axes - see Table I.) The -- mode can be thought of as a dipole mode in the transverse direction. As a result it shows a strong frequency dependence on the transverse dimension w . When $w/L = 1$, the -- mode is identical to the ++ mode rotated spatially 90 degrees. Consequently, the two trajectories coalesce as $w/L \rightarrow 1$, when the zoning is 5x5 so as to preserve symmetry in the numerical mode. The failure of the 5x3 zone case is due to the reasons outlined in the previous section.

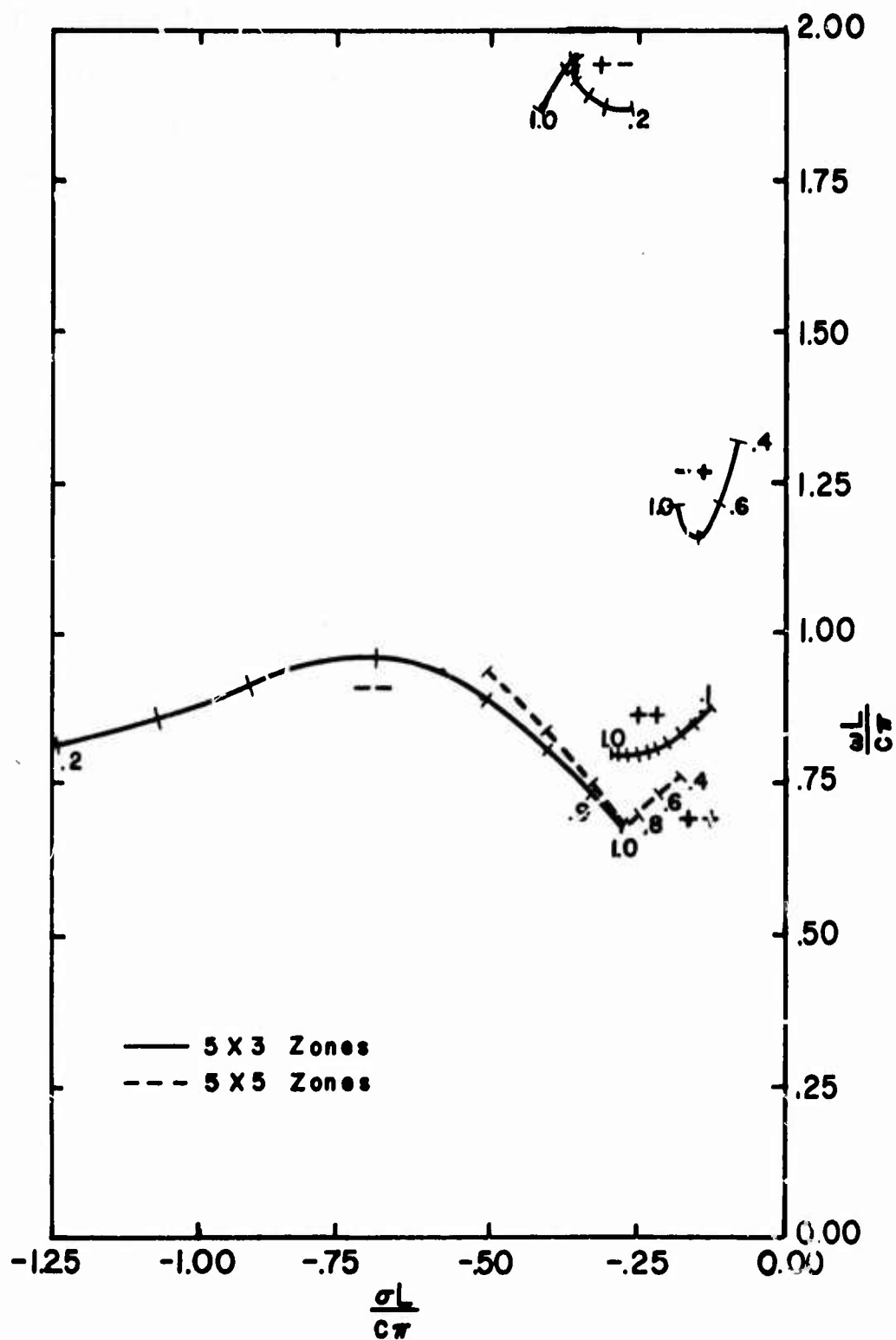


Figure 9. Pole Trajectories as Computed with Zoning Fixed

SECTION VIII

CONCLUSIONS

The application of SEM to the equivalent problems of the perfectly conducting rectangular plate and the rectangular aperture in a perfectly conducting screen has been conducted with partial success. The applicability of SEM and the computational feasibility of determining SEM quantities are demonstrated. It is further demonstrated that by careful program construction, the computational costs of numerical treatment of two-dimensional problems can be made quite reasonable. The cost of generating a matrix and calculating its determinant by the methods discussed herein is less than 10 cents for each value of s .

Difficulties arise in the subsectionally constant current zoning because of a failure to properly model the edge conditions. Whereas Rahmat-Samii and Mittra (ref. 3) obtained good radar cross-section results with such a zoning scheme, the pole locations are highly sensitive to the zoning. Radar cross-section is a far-field quantity, and the integrated effects of the errors are small there. The pole locations, on the other hand, depend strongly on the dimensions of the structure, and it is evident that the plate size must be brought to bear in a precise fashion for the pole locations to be correct.

The edge problem can be remedied by imposing the edge conditions (3) directly in the basis set elements for edge zones. Davis has done this successfully for the circumferential component of current on a thick cylindrical scatterer (ref. 8). The circumferential current

-
8. Davis, W. A., "Numerical Solutions to the Problem of Electromagnetic Radiation and Scattering by a Finite Cylinder," Ph.D. Thesis, University of Illinois, 1974.

component is singular at the ends of the cylinder. The introduction of the singular basis element will produce a significant complication to the problem in that a second singularity, that of the current, must be integrated analytically in order to implement the model with edge conditions imposed. An independent check on program accuracy is dictated for a problem of this scope before proceeding with the edge condition approach.

APPENDIX A

THE SELF-PATCH INTEGRATION

The term for the interaction matrix for $IDIF = JDIF = 0$, i.e., where the match point lies at the center of the source patch, can be written

$$I_s = 4 \int_0^{\Delta x/2} \int_0^{\Delta y/2} K(0,0|x',y') dx' dy' \quad (A1)$$

This presumes a unit amplitude expansion pulse over the patch whose dimensions are Δx and Δy . The symmetry of the kernel with respect to both x' and y' is employed in the forming of (A1). This integral can be transformed to polar coordinates as

$$\begin{aligned} I_s &= 4 \left\{ \int_{\phi=0}^{\tan^{-1} \frac{\Delta y}{\Delta x}} \int_{\rho=0}^{\frac{\Delta x}{2 \cos \phi}} \exp[-s\rho/c] d\rho d\phi \right. \\ &\quad \left. + \int_{\phi=\tan^{-1} \frac{\Delta y}{\Delta x}}^{\pi/2} \int_{\rho=0}^{\frac{\Delta x}{2 \cos \phi}} \exp[-s\rho/c] d\rho d\phi \right\} \\ &= -\frac{4c}{s} \left\{ \int_{\phi=0}^{\tan^{-1} \frac{\Delta y}{\Delta x}} [\exp(-s\Delta x \sec \phi/2c) - 1] d\phi \right. \\ &\quad \left. + \int_{\phi=\tan^{-1} \frac{\Delta y}{\Delta x}}^{\pi/2} [\exp(-s\Delta y \csc \phi/2c) - 1] d\phi \right\} \quad (A2) \end{aligned}$$

If the exponential functions in the integrand are then expanded in McLaurin series, the resulting terms of powers of secants and cosecants possess tabulated integrals. The result for three terms retained in the series is

$$\begin{aligned}
I_s \approx & -\frac{4c}{s} \left\{ -\frac{s\Delta x}{2c} \cdot \frac{1}{2} \ln q_y + \frac{1}{2} \left(\frac{s\Delta x}{2c} \right)^2 \frac{\Delta y}{\Delta x} \right. \\
& - \frac{1}{6} \left(\frac{s\Delta x}{2c} \right)^3 \frac{\Delta x(\Delta x^2 + \Delta y^2)^{1/2}}{2\Delta y^2} - \frac{s\Delta y}{2c} \cdot \frac{1}{2} \ln q_x \\
& \left. + \frac{1}{2} \left(\frac{s\Delta y}{2c} \right)^2 \frac{\Delta x}{\Delta y} - \frac{1}{6} \left(\frac{s\Delta y}{2c} \right)^3 \frac{\Delta y(\Delta x^2 + \Delta y^2)^{1/2}}{2\Delta y^2} \right\} \quad (A3)
\end{aligned}$$

where

$$q \begin{pmatrix} x \\ y \end{pmatrix} = \frac{[(\Delta x^2 + \Delta y^2)^{1/2} + (\Delta x)]}{[(\Delta x^2 + \Delta y^2)^{1/2} - (\Delta x)]}$$

APPENDIX B

THE SPARSE MATRIX ALGORITHMS

1. Introduction

This Appendix describes the algorithmic approach to minimize the computation time involved in Gaussian elimination triangularization of systems of matrix equations which are "sparsely coupled." The term "sparsely coupled" as applied in this Appendix implies the matrix equation form given in (B1).

$$[M] [X] = \begin{bmatrix} M_1 & 0 \\ 0 & M_2 \end{bmatrix} \begin{bmatrix} X \\ X \end{bmatrix} = \begin{bmatrix} B \\ B \end{bmatrix} \quad (B1)$$

It is clear that in this form the sole coupling between the "upper" and "lower" systems of equations is contained in the matrix M_2 . Generally, the number of columns in M_2 is small compared with the order of the overall system.

An algebraic approach to exploiting the sparceness of (B1) results in solutions which are given in terms of the several submatrices and their inverses. (See, for example, ref. 9.) It is well-known, however, that it is sufficient for the purposes of determinant calculation and equation solution to triangularize the composite matrix in (B1). The triangularization process involves fewer operations than the diagonalization necessary for the calculation of an inverse.

-
9. Dunaway, O. C., "A Numerical Solution for the Distribution of Time-Harmonic Electromagnetic Fields in an Arbitrary Shaped Aperture in a Ground Screen," M.S. Thesis, University of Mississippi, 1974.

$$\begin{bmatrix} \begin{bmatrix} \text{CMAT1} \\ (\text{NI1} \times \text{NJ1}) \end{bmatrix} & \begin{bmatrix} 0 \\ (\text{NI1} \times \text{NJ2}) \end{bmatrix} & \begin{bmatrix} \text{CMAT3} \\ (\text{NI3} \times \text{NJ3}) \end{bmatrix} \\ \begin{bmatrix} 0 \\ (\text{NI2} \times \text{NJ1}) \end{bmatrix} & \begin{bmatrix} \text{CMAT2} \\ (\text{NI2} \times \text{NJ2}) \end{bmatrix} & \end{bmatrix}$$

$$\text{NI3} = \text{NI1} + \text{NI2}$$

$$\text{NJ3} = \text{NI3} - \text{NJ1} - \text{NJ2}$$

(a)

$$\begin{bmatrix} \begin{bmatrix} \neq 0 \\ 0 \\ 0 \end{bmatrix} & \begin{bmatrix} 0 \\ \text{CMAT4} \\ 0 \end{bmatrix} & \begin{bmatrix} \neq 0 \\ \neq 0 \\ 0 \end{bmatrix} \\ \begin{bmatrix} \text{CMAT4} \\ (\text{NI4} \times \text{NI2}) \end{bmatrix} & \begin{bmatrix} 0 \\ 0 \end{bmatrix} & \end{bmatrix}$$

$$\text{NI4} = \text{MAX} (\text{NI1} - \text{NJ1}, 0)$$

(b)

Figure B1. Submatrix Organization for the Sparse Matrix Algorithms, a) the Original Matrix, and b) Triangularized Form with the Generated CMAT4

This Appendix describes an algorithmic exploitation of the sparseness of the composite matrix in (B1). That is, a numerical process is given whereby only the non-zero subelements are stored and operated on, with the computational operations being those which render the composite matrix M upper triangular. The determinant of the composite matrix results directly from this triangularization. A solution for X in (B1) requires a backsolving process involving the triangularized form of M and a vector resulting from applying the elimination operations to the vector B . Routines to perform these operations are given as well.

Appendix C gives listings of the routines built on this algorithm. The triangularization routine is named SPRHOM. The backsolving procedure is performed by the entry HOMSLV to the routine SPRSLV. (The name entry SPRSLV backsolves an inhomogeneous system and is not used for present purposes.)

2. The Algorithm

The routine SPRHOM is simply an implementation of the Gaussian elimination procedure with maximum pivot selection applied to the composite matrix M in (B1) viewed as a whole. The sparseness of M is exploited by storing only the non-zero submatrices in (B1) and skipping the operations involving zero elements. The result is a substantial saving in both storage and computation time.

The forms of the input and of the end product for the triangularization are given in Figure (B1) with the sizes of the respective submatrices defined. It is recalled that the fundamental process in the Gaussian elimination procedure is the elimination of all or part of the elements of a column of a matrix with respect to a pivot element, commonly the element of greatest magnitude in the column. That is, for a column under process, the row

containing the main diagonal element of the matrix which falls in that column. All or part of the elements not on the main diagonal are "eliminated" or made zero by subtraction of some multiple of the row containing the column maximum. In the triangularization procedure, the part of the column comprising elements lying below the main diagonal after row exchange are eliminated. If the matrix is a part of a system of equations with non-zero right-hand side, the row operations of exchange and subtraction of a constant multiple of another row must be performed on the corresponding elements of the right-hand side vector as well.

The algorithm of the routine SPRHOM operates according to the Gaussian elimination procedure described above. However, the three submatrices CMAT1, CMAT2, and CMAT3 are stored individually. In addition, the routine generates a submatrix CMAT4 in the course of selecting pivots for the columns contained in CMAT2. Further, the "elimination" of elements of submatrices that are zero a priori is skipped. The result is significant storage and time economy.

In order to minimize logic decisions in the routine, it is organized to operate sequentially through the partitioned matrix. The steps are as follows (it is convenient to follow the thinking of these steps by tracing the location diagonal of the composite with the aid of Table B1):

- a. Perform conventional Gaussian elimination to zero the elements CMAT1(I,J) for $I > J$, i.e., the elements below the main diagonal of M. Choose maximum pivot elements in conventional fashion. Carry row operations into CMAT3.
- b. Create CMAT4 by row swapping with CMAT2 so as to choose maximum pivot elements. Perform elimination to zero CMAT4 elements for $I > J$ and the entire column of CMAT2. The number of rows created in CMAT2 is $NI4 = NI1 - NJ1$, the amount by which CMAT1 is over-square. Carry row operations across into CMAT3.

- c. Choose maximum pivot rows in columns of CMAT2 with $J > NI4$ and swap with rows given by $I = J - NI4$ (the rows containing the Jth column diagonal element of the composite). Conduct elimination to zero elements with $I > J + NI4$. Carry row operations into CMAT3.
- d. Conduct conventional pivot selection and elimination on CMAT3 to zero elements CMAT3(I, J) with $I > J + NJ1 + NJ2$.

In each column elimination operation, the pivot value is multiplied into a product accumulator to produce a value for the determinant of the composite matrix. The sign of this product is changed at each row swap in accord with the rules of matrix algebra row operations.

The backsolving routine SPRSLV with its entry HOMSLV operate in a straightforward manner on the submatrices as reduced by SPRHOM. Details are omitted here, but the routines may be easily followed in a row-by-row flow from the bottom of the composite matrix, if one keeps in mind the row index relations of column 4 of Table B1. The entry HOMSLV assumes a zero determinant value resulted (approximately) from SPRHOM and the last element of the solution vector is picked as unity. (The zero determinant results from a zero falling at the last diagonal location in maximum pivoting Gaussian elimination.) The remainder of the homogeneous solution vector is backsolved conventionally relative to this last element. The vector is then renormalized so that its maximum element is unity.

Table B1

PRIMARY INDEXING QUANTITIES IN THE ALGORITHM

Submatrix	Size of ¹ Submatrix	Indices of Main ¹ Diag. of Compos.	Relative Row ² Index of CMAT3 and CRHS
CMAT1	NI1 x NI2	(J,J)	I3 = I
CMAT4	NI1 - NJ1 x NJ2 (can be null)	(J,J)	I3 = I + NJ1
CMAT2	NI2 x NJ2	(J - (NI1 - NJ1), J)	I3 = I + NI1
CMAT3	NI1 + NI2 x NI1 + NI2 - NJ1 - NJ2	(J + NJ1 + NJ2, J)	I3 = I3

1. Quantities given in terms of input parms. to the routine. Related internal quantities are given in Figure B1.
2. Relative to I, the row index of the submatrix in question.

APPENDIX C
PROGRAM LISTINGS

All code compileable on IBM OS/360 and OS/370 FORTRAN levels G or H. The routine ZANLYT and its service routine UERTST is taken from the program library FORTUOI made available by the Computer Services Office, University of Illinois at Urbana-Champaign, Urbana, Illinois 61801. The routines BSLJZ and BSCJZ are taken from the International Mathematical and Statistical Library (IMSL). They may not be reproduced apart from this application program package. The IMSL library is available by subscription from IMSL, Inc., 6100 Hillcroft, Suite 510, Houston, Texas 77036.

```

* POLE SEARCH PROGRAM FOR S E M FORMULATION OF THIN-PLATE SCATTERER 00010
C BY L W PEARSON 8/74 00020
C 00030
IMPLICIT REAL*8(A,B,D-H,C-Z),COMPLEX*16(C) 00040
COMMON /GE3M/ XSYM,YSYM,W,NX,NY,IPREAS(10),JPREAS(10),NPPEI,NPREJ 00050
INTEGER MES(4,2)/'SYMM','ETRI','C',' ','ANTI','SYMM','ETRI','C'/ 00060
DATA C/(3.008,0.00)/,PLUS/'+'/,PI/3.141592653589793/ 00070
DATA HX/'X'/,HY/'Y'/ 00080
EXTERNAL CPLATE 00090
DIMENSION CX(20),INFFR(20) 00100
LOGICAL LAUTO 00110
100 READ(5,1,END=999) XSYM,YSYM,NX,NY,W0,WS,W0,CSUNCP,LAUTO 00120
1 FORMAT(2A1,2X,2I3,5F10.4,T80,L1) 00130
IMX=1 00140
IMY=1 00150
IF(XSYM.NE.PLUS) IMX=2 00160
IF(YSYM.NE.PLUS) IMY=2 00170
VW=(WM-W0)/WS 00180
IF(NW.GT.0) GO TO 105 00190
NW=-NW 00200
WS=-WS 00210
105 IF(WS*NW.LT.WM-W0) NW=NW+1 00220
DO 200 IW=1,NW 00230
W=W0+(IW-1)*WS 00240
IF(.NOT.LAUTO) GO TO 140 00250
                                SKIP PAST AUTO ZONING 00260
C 00270
C ROUTINE TO DETERMINE NO OF EXPANSION PULSES BASED ON ELECTRICAL 00280
C SIZE OF PLATE 00290
C 00300
TESTWV=.1885010/DABS(DIMAG(CSUNCP)) 00310
C 00320
C ////////// 00330
C 00340
NPPWVL=6 00350
C 00360
C ////////// 00370
C 00380
FLENX=1/TESTWV 00390
NX=IDINT(FLENX*NPPWVL) 00400
IF(DFLOAT(NX).LT.FLENX*NPPWVL) NX=NX+1 00410
FLENY=W/TESTWV 00420
NY=IDINT(FLENY*NPPWVL) 00430
IF(DFLOAT(NY).LT.FLENY*NPPWVL) NY=NY+1 00440
NX=MIN0(NX,5) 00450
NY=MIN0(NY,5) 00460
C 00470
C BEGIN SETUP FOR ALTERNATE EDGE PATCH REASSIGNMENT 00480
C 00490
140 NPPEI=(NX+2)/3 00500
NPREJ=(NY+2)/3 00510
IF(NX-2*NPPEI+2.LE.1.AND.NPPEI.GT.1) NPPEI=NPPEI-1 00520
IF(NY-2*NPREJ+2.LE.1.AND.NPREJ.GT.1) NPREJ=NPREJ-1 00530
DO 110 I=1,NPPEI 00540
IPREAS(NPPEI+1-I)=NX-3*I+3 00550
110 CONTINUE 00560
DO 120 J=1,NPREJ 00570
JPREAS(NPREJ+1-J)=NY-3*J+3 00580
120 CONTINUE 00590
C LOCATIONS WHERE X-DIRECTED CURRENT 00600
C IS SET TO ZERO GIVEN BY SUBS01 00610

```


C	PTS (IX,JPREAS) AND Y-DIRECTED BY	00620
C	(IPFEAS,NY)	00630
C		00640
	WRITE(6,2) W,CSUNDP	00650
2	FORMAT('ENTER ITERATION',/, 'OSHAPE RATIO =',F5.3,5X,	00660
	1'STARTING FREQ =',2D12.4)	00670
	WRITE(6,3)	00680
3	FORMAT('O',10X,'CUR SYMMETRY',6X,'PULSES',3X,'PREASSIGN LOC' 'NS')	00690
	WRITE(6,4) HX,(MES(I,IMX),I=1,4),NX,(IPREAS(J),J=1,NPREI)	00700
4	FORMAT(' ',A1,'-DIR',5X,4A4,16,5X,1013)	00710
	WRITE(6,4) HY,(MES(I,IMY),I=1,4),NY,(JPREAS(J),J=1,NPPEJ)	00720
	WRITE(6,5)	00730
5	FORMAT('O',11X,'COMPLEX FREQ',17X,'DETERMINANT')	00740
	CX(1)=CSUNDP	00750
	CALL ZANLYT(CPLATE,1.0-50,4,0,1,1,CX,100,INFER,IFF)	00760
	WRITE(6,6) CX(1)	00770
6	FORMAT('RETURN FROM ITERATION',/, 'OPCLE LOC' 'N',2E12.4)	00780
	CALL MODE	00790
	CSUNDP=CX(1)	00800
200	CONTINUE	00810
	GO TO 100	00820
999	STOP	00830
	END	00840

Reproduced from
best available copy.

SUBROUTINE MODE	0085C
IMPLICIT REAL*8(A,B,D-H,O-Z),COMPLEX*16(C)	00860
COMMON /MAT/ CMATX(25,25),CMATY(25,25),CHOM(50,10),CMAT4(10,25),	0087C
1 NPTCHS,NDIM1,NDIMCI,NDIMCJ,NORD	0088C
COMMON /GEM/ XSYM,YSYM,W,NX,NY,IPREAS(10),JPPEAS(10),NPRED,NPREJ	00890
DIMENSION CPRX(5,5),CPRY(5,5)	00900
DIMENSION CJ(50)	00910
NPRED=NPRED+NPRED	00920
NPRED=NPRED-1	00930
NPRED=NPRED-1	00940
CALL HOMSLV(CMATX,NPTCHS,NPTCHS-NPREJ,NDIM1,NDIM1,	00950
1 CMATY,NPTCHS,NPTCHS-NPREI,NDIM1,NDIM1,	00960
2 CHOM,NDIMCI,NDIMCJ,CMAT4,NDIMCJ,NDIM1,CJ,NORD)	00970
NXM1=NX-1	00980
NYM1=NY-1	00990
NSURS=0	01000
DO 470 JS=1,NY	01010
DO 450 IS=1,NXM1	01020
J=(JS-1)*NX+IS	01030
CPRX(IS,JS)=CJ(J-NSURS)	01040
JM=J-NSURS	01050
450 CONTINUE	01060
J=JS*NX	01070
IF(JS.NE.JPREAS(NSURS+1)) GO TO 460	01080
NSURS=MINO(NSURS+1,NPREJM)	01090
CPRX(NX,JS)=(0.,0.)	01100
GO TO 470	01110
460 CPRX(NX,JS)=CJ(J-NSURS)	01120
470 CONTINUE	01130
DO 500 IS=1,NX	01140
DO 500 JS=1,NYM1	01150
J=(JS-1)*NX+IS	01160
CPRY(IS,JS)=CJ(NPTCHS-NPREJ+J)	01170
500 CONTINUE	01180
NSURS=0	01190
DO 530 IS=1,NX	01200
J=NYM1*NX+IS	01210
IF(JS.NE.JPREAS(NSURS+1)) GO TO 510	01220
CPRY(IS,NY)=(0.,0.)	01230
NSURS=MINO(NSURS+1,NPREIM)	01240
GO TO 530	01250
510 CCRY(IS,NY)=CJ(NPTCHS-NPREJ+J-NSURS)	01260
530 CONTINUE	01270
WRITE(6,16)	01280
16 FORMAT('O*****NATURAL MODE*****',/, 'OX-DIRECTED CURRENT')	01290
DO 540 I=1,NX	01300
WRITE(6,17) (CPRX(I,J),J=1,NY)	01310
17 FORMAT('O',5(2D12.4,2X))	01320
540 CONTINUE	01330
WRITE(6,18)	01340
18 FORMAT('OY-DIRECTED CURRENT')	01350
DO 550 I=1,NX	01360
WRITE(6,17) (CPRY(I,J),J=1,NY)	01370
550 CONTINUE	01380
WRITE(6,19)	01390
19 FORMAT('OHOMOGENEOUS EXPANSION COEFF'S')	01400
WRITE(6,17) (CJ(2*NPTCHS-NPRE+I),I=1,NPRE)	01410
RETURN	01420
END	01430

C	COMPLEX FUNCTION CPLATE*16(CSUNOP)	01440
C	DETERMINANT EVALUATION ROUTINE FOR HALLEN-TYPE AUGMENTED MOMENT	01450
C	MATRIX FOR THE THIN PLATE SCATTERER	01460
C	FOR S E M APPLICATIONS	01470
C	BY L W PEARSON 8/74	01480
C		01490
	IMPLICIT COMPLEX*16(C),REAL*8(A,R,D-H,C-Z)	01500
	COMMON /GEO/ XSYM,YSYM,W,NX,NY,IPREAS(10),JPREAS(10),NPRI,NPREJ	01510
	COMMON /MAT/ CMATX(25,25),CMATY(25,25),CHCM(50,10),CMAT4(10,25),	01520
	INPTCHS,NDIM1,NDIMCT,NDIMCJ,NORD	01530
	REAL*8 DRARG,DIMARG,DBRES(20),DIMBFS(20),DUM1(20),DUM2(20),DUM3(20	01540
	1),DUM4(20)	01550
	DIMENSION CINTER(10,10),CINTX(25),CINTY(25),CCOSTM(10),CSINTM(10)	01560
	INTEGER MES(4,2)/'SYMM','ETRI','C',' ','ANTI','SYMM','ETRI','C'/	01570
	DATA C/(3.D08,0.D0)/,PLUS/'+'/,PI/3.141592653589793/	01580
	NDIM1=25	01590
	NDIMCT=50	01600
	NDIMCJ=10	01610
	NDIM=50	01620
C		01630
C	FORMULATION SETUP ROUTINES	01640
C		01650
	TMX=1	01660
	TMV=1	01670
	IF(XSYM.EQ.PLUS) TMX=2	01680
	IF(YSYM.EQ.PLUS) TMV=2	01690
C		01700
C	IM(X/Y)=2 INDICATES ANTISYMMETRIC	01710
C	DISTR OF X-DIRECTED CURRENT WRT X	01720
	/Y AXIS	01730
C	NPTCHS=NX*NY	01740
		01750
C		01760
	NXM1=NX-1	01770
	NVM1=NY-1	01780
	EDGFAC=0.5	01790
C		01800
	EDG2=EDGFAC*EDGFAC	01810
C		01820
	DX=1./((FLOAT(2*NX-2)+2*EDGFAC)	01830
	DY=W/((FLOAT(2*NY-2)+2*EDGFAC)	01840
	NXT2=NX*2	01850
	NYT2=NY*2	01860
	CS=CSUNOP/2/	01870
C		01880
	INPTCHS=13	01890
	DXINT=DX/12	01900
	DYINT=DY/12	01910
C		01920
	NSYMX=(-1)**TMX	01930
	NSYMY=(-1)**TMV	01940
C		01950
	NSMTY=NSYMY	01960
	NSMTI=NSYMX*NSYMY	01970
	NSMTV=NSYMX	01980
C		01990
C		02000
	NINDX=2	02010
	IF(NSMTI.GT.0) NINDX=1	02020
C		02030
	NSCCS=1	02040
	IF(XSYM.EQ.PLUS) NSCCS=2	

C		VSCOS = 2 INDICATES EVEN SYMM WRT	02050
C		Y FOR X DIR CUAR (I E COSINE EXPA	02060
C		NSION OF HOMOGENEOUS SOLID)	02070
C			02080
C	NPRE=NPRI+NPRIJ		02090
C		TOT NO OF PREASSIGNED CURR VALHS	02100
C	NPRIJM=NPRIJ-1		02110
C	NPRIIM=NPRI-1		02120
C	NPRIJ=NPRI+1		02130
C	END OF INPUT/SPECIFICATION ROUTINES		02140
C			02150
C	ROUTINE TO FILL INTERACTION MATRIX FROM WHICH MOMENT MATRIX IS		02160
C	CONSTRUCTED		02170
C			02180
C	DIAG=DSQRT(DX*DX+DY*DY)		02190
C	ALNXTM=2*DLOG((DIAG+DY)/DX)		02200
C	ALNYTM=2*DLOG((DIAG+DX)/DY)		02210
C	DYBDX=DY/DX		02220
C	DXBDY=DX/DY		02230
C	CSDX=CS*DX		02240
C	CSDY=CS*DY		02250
C	CSDX2=CSDX*CSDX		02260
C	CSDX3=CSDX*CSDX2		02270
C	CSDX4=CSDX2*CSDX2		02280
C	CSDY2=CSDY*CSDY		02290
C	CSDY3=CSDY*CSDY2		02300
C	CSDY4=CSDY2*CSDY2		02310
C		COMPONENT TERMS FOR SELF-PATCH SE	02320
C		RIES	02330
C	CXTERM=-0.500*CSDX*ALNXTM+0.500*CSDX2*DYBDX-CSDX3*(DXBDY*DIAG/(12*		02340
C	1DY)+ ALNXTM/24)+CSDX4*DYBDX*(1+DYBDX*DYBDX/3)/24		02350
C	CYTERM=-0.500*CSDY*ALNYTM+0.500*CSDY2*DXBDY-CSDY3*(DYBDX*DIAG/(12*		02360
C	1DX)+ ALNYTM/24)+CSDY4*DXBDY*(1+DXBDY*DXBDY/3)/24		02370
C		CALC INDIV SERIES FOR SELF-INTER	02380
C	CINTER(1,1)=-2/CS*(CXTERM+CYTERM)		02390
C		COMPUTE SELF-INTERACTION	02400
C	DO 220 IS=1,2		02410
C	XS=(FLOAT(IS)-1.500)*DX		02420
C	DO 220 JS=1,2		02430
C		LOOP TO CALC ADJACENT PATCH INTER	02440
C	IF(JS+JS.EQ.1) GO TO 220		02450
C	YS=(FLOAT(JS)-1.500)*DY		02460
C	DO 210 JINT=1,INTPTS		02470
C	XP=FLOAT(JINT-1)*DXINT		02480
C		NUMER INT WRT X LOOP	02490
C	X2TM2=XS+XP		02500
C	X2TM2=X2TM2*X2TM2		02510
C	DO 200 JINT=1,INTPTS		02520
C	YP=FLOAT(JINT-1)*DYINT		02530
C		NUMER INT WRT Y LOOP	02540
C	Y2TM=YS+YP		02550
C	P=DSQRT(X2TM2+Y2TM*Y2TM)		02560
C	CINTY(JINT)=CDEXP(-2*CS*P)/P		02570
C		EVAL INTEGRAND	02580
C	200 CONTINUE		02590
C	CALL DWEDDL(CINTY,INTPTS,DYINT,CINTX(IINT))		02600
C		INT WRT Y TO YIELD X INTEGRAND	02610
C	210 CONTINUE		02620
C	CALL DWEDDL(CINTX,INTPTS,DXINT,CINTER(IS,JS))		02630
C		INT WRT X	02640
C	220 CONTINUE		02650

	DO 250 JS=1,NXT2	02660
	X2TM2=DFLOAT(JS-1)*DX	02670
	X2TM2=X2TM2*X2TM2	02680
	DO 250 JS=1,NYT2	02690
C		LOOPS FOR REMAINDER OF INTERACTION
C		N CALC'D BY ONE-PT RECTANG RULE
	IF (IS+JS.LT.4.OR.IS.EQ.2.AND.JS.EQ.2) GO TO 250	02710
	Y2TM=FLOAT(JS-1)*DY	02720
	R=DSORT(X2TM2+Y2TM*Y2TM)	02730
	CINTER(JS,JS)=CDEXP(-2*CS*R)/R*DX*DY	02740
250	CONTINUE	02750
C	END OF LOOP TO FILL INTERACTION MATRIX	02760
C		02770
C	BEGIN CONSTRUCTION OF MOMENT MATRIX	02780
C		02790
	DO 350 IM=1,NX	02800
	DO 350 JM=1,NY	02810
C		02820
C		INDEXING OF MATCH-POINTS OVER ENT
C		IRE QUADRANT
	I=(JM-1)*NX+IM	02830
C		02840
C		ONE-DIM MATCH-PT INDEX GEN'D
C		COL'WISE ALONG X-DIRECTION
	NSURS=0	02850
	DO 330 JS=1,NYM1	02860
C		02870
C		INDEX OVER SOURCE PATCHES Y-DIR
	JD1=IABS(JM-JS)+1	02880
C		02890
C		1ST AND 2ND QUAD J 'DIFFERENCE
C		INDEX'
	JD2=JM+JS	02900
C		02910
C		3RD & 4TH QUAD J 'DIFFERENCE
C		INDEX'
C		02920
C		NOTE THAT 'DIFFERENCE INDICES' AR
C		E = 'INDEX DIFFERENCE' +1 FOR THE
C		SAKE OF FORTRAN INDEXING
	DO 310 IS=1,NXM1	02930
C		02940
C		INDEX OVER SOURCE PATCHES X-DIR
	ID1=IABS(IS-IM)+1	02950
C		02960
C		1ST & 4TH QUAD I 'DIFF INDEX'
C		ID2=IS+IM
C		02970
C		2ND & 3RD QUAD I 'DIFF INDEX'
	J=(JS-1)*NX+IS	02980
C		02990
C		ONE-DIM SOURCE-PT INDEX
	CO=CINTER(ID1,JD1)+NSMII*CINTER(ID2,JD2)	03000
C		03010
C		SUM OF SOURCE CONT FROM QI & QIII
	CE=NSMII*CINTER(ID2,JD1)+NSMIV*CINTER(ID1,JD2)	03020
C		03030
C		SUM OF SOURCE CONT FROM QII & QIV
	CMATX(I,J-NSURS)=CO+CE	03040
C		03050
C		SURMAT ENTRY FOR X-DIR CURR'S
	CMATY(I,J)=CO-CE	03060
C		03070
C		SURMAT ENTRY FOR Y-DIR CURR'S
C		03080
C		NOTE THAT EVEN Q'S CONT NEGATIVE
C		FOR Y-DIR CURR'S
310	CONTINUE	03090
C		03100
C	END OF SOURCE LOOP FOR INTERIOR PATCHES	03110
C		03120
C	CONSTRUCTION OF SOURCE TERMS FROM ABS(X)=A EDGE FOLLOWS	03130
C		03140
	ID1=IABS(NX-IM)+1	03150
	ID2=NX+IM	03160
	J=JS*NX	03170
		03180
		03190
		03200
		03210
		03220
		03230
		03240
		03250
		03260

```

C      CD=CENTER(ID1,JD1)+NSMIII*CENTER(ID2,JD2)                                03270
C      SUM OF SOURCE CONT FROM QI & QIII                                         03280
CE=NSMII*CENTER(ID2,JD1)+NSMIV*CENTER(ID1,JD2)                                03290
C      SUM OF SOURCE CONT FROM QII & QIV                                         03300
CMATY(I,J)=(CD-CE)*EDGEFAC                                                    03310
C      SURMAT ENTRY FOR Y-DIR CURR'S                                           03320
C      NOTE THAT EVEN Q'S CONT NEGATIVE                                         03330
C      FOR Y-DIR CURR'S                                                         03340
      IF(JS.NE.JPREAS(NSURS+1)) GO TO 325                                         03350
      NSURS=MING(NSURS+1,NPREJM)                                                 03360
      GO TO 330                                                                    03370
325  CMATX(I,J-NSURS)=(CE+CD)*EDGEFAC                                           03380
C      SURMAT ENTRY FOR X-DIR CURR'S                                           03390
C      END ROUTINE FOR ABS(X)=R EDGE TERMS                                       03400
C      03410
C      03420
330  CONTINUE                                                                    03430
C      03440
C      END LOOP OVER Y COORD FOR INTERIOR PATCHES                             03450
C      03460
C      BEGIN ROUTINE FOR CONSTRUCTION OF SOURCE TERMS FOR ABS(Y)=R EDGE        03470
C      03480
      JD1=IABS(NY-JM)+1                                                           03490
      JD2=NY+JM                                                                    03500
      NSURSJ=NSURS                                                                03510
      NSURS=0                                                                      03520
      DO 340 IS=1,NXM1                                                            03530
C      INDEX DOWN X COORD'S INTERIOR PATCHES                                  03540
C      03550
      ID1=IABS(IS-IM)+1                                                           03560
      ID2=IS+IM                                                                    03570
      J=(NYM1)*NX+IS                                                             03580
      CD=CENTER(ID1,JD1)+NSMIII*CENTER(ID2,JD2)                                03590
C      SUM OF SOURCE CONT FROM QI & QIII                                         03600
      CE=NSMII*CENTER(ID2,JD1)+NSMIV*CENTER(ID1,JD2)                            03610
C      SUM OF SOURCE CONT FROM QII & QIV                                         03620
      CMATX(I,J-NSURSJ)=(CE+CD)*EDGEFAC                                           03630
C      SURMAT ENTRY FOR X-DIR CURR'S                                           03640
      IF(IS.NE.IPREAS(NSURS+1)) GO TO 335                                         03650
      NSURS=MING(NSURS+1,NPREJM)                                                 03660
      GO TO 340                                                                    03670
335  CMATY(I,J-NSURS)=(CD-CE)*EDGEFAC                                           03680
C      SURMAT ENTRY FOR Y-DIR CURR'S                                           03690
C      NOTE THAT EVEN Q'S CONT NEGATIVE                                         03700
C      FOR Y-DIR CURR'S                                                         03710
340  CONTINUE                                                                    03720
C      03730
C      END ROUTINE FOR ABS(Y)=R EDGE                                             03740
C      03750
C      CONSTRUCTION OF CORNER SOURCE TERM                                       03760
C      03770
      ID1=IABS(NX-IM)+1                                                           03780
      ID2=NX+IM                                                                    03790
      J=NX*NY                                                                    03800
      CD=CENTER(ID1,JD1)+NSMIII*CENTER(ID2,JD2)                                03810
C      SUM OF SOURCE CONT FROM QI & QIII                                         03820
      CE=NSMII*CENTER(ID2,JD1)+NSMIV*CENTER(ID1,JD2)                            03830
C      SUM OF SOURCE CONT FROM QII & QIV                                         03840
      IF(NY.NE.JPREAS(NPREJ)) CMATX(I,J-NPREJM)=(CE+CD)*EDG2                   03850
C      SURMAT ENTRY FOR X-DIR CURR'S                                           03860
      IF(NX.NE.IPREAS(NPREI)) CMATY(I,J-NPREIM)=(CD-CE)*EDG2                   03870

```

C		SURMAT ENTRY FOR Y-DIR CURVES	04380
C		NOTE THAT EVEN Q'S CAN BE NEGATIVE	04390
C		FOR Y-DIR CURVES	04390
C	350	CONTINUE	043910
C		END OF MOMENT MATRIX INTERACTION CONSTRUCTION	043920
C		BEGIN ROUTINE TO ENTER HOMOGENEOUS SOLIN EXPANSION COEFFS IN MATRIX	043930
C			043940
C	360	NRES=2*NPPE	043950
C		HIGHEST ORDER BESSEL FUNCTION IN	043960
C		HOMOGENEOUS SOLIN EXPANSION	043970
C		IF (NINDEX.EQ.2) NRES=NRES-1	043980
C		ONE LESS IF EVEN INDEX EXPANSION	043990
C			044000
C		DO 400 IM=1,NX	044010
C		X=(FLOAT(IM)-0.500)*DX	044020
C		DO 400 JM=1,NY	044030
C		Y=(FLOAT(JM)-0.500)*DY	044040
C		I=(JM-1)*NX+IM	044050
C		INDEXING THRU MATCH-PTS	044060
C		PHI=DATAN(Y/X)	044070
C		RHO=DSQRT(X*X+Y*Y)	044080
C		POLAR COORDS OF MATCH-PTS	044090
C		DRARG=2*DIMAG(CS)*RHO	044100
C		DMARG=-2*DREAL(CS)*RHO	044110
C		ARGUMENT OF BESSEL FN'S	044120
C		IF (DABS(DMARG/DRARG).LT.1.E-20) GO TO 364	044130
C		IF REAL ARG SKIP TO REAL BES CALL	044140
C		CALL BSCJZ(DRARG,DMARG,DRRES,DMRES,NRES,0.00,16,IFR,DUM1,DUM2,D	044150
C		IUM3,DUM4)	044160
C		GET TABLE OF BESSEL FUNCTIONS	044170
C		GO TO 368	044180
C	364	CALL BSLJZ(DRARG,DRRES,NRES,0.00,16,IFR,DUM1,DUM2)	044190
C		CALL ZEROZ(DMRES,2*(NRES+1))	044200
C		CALL ZEROZ(DMRES,2*(NRES+1))	044210
C		SET UP PURE REAL RES FUNCTIONS	044220
C	368	CCOSTM(1)=0	044230
C		CSINTM(1)=0	044240
C		ZERO 1ST TERM COEF CONSTRUCTION	044250
C		VECTORS	044260
C		DO 370 II=1,NPREP1	044270
C		INDEX THRU CALC OF COEF CONSTR	044280
C		VECTOR	044290
C		INDEX=2*II-NINDEX	044300
C		CALC SERIES INDEX	044310
C		IF (INDEX.EQ.0) GO TO 370	044320
C		SKIP CALC OF BELOW TERM FOR ZERO	044330
C		INDEX - IT WAS SET TO ZERO ABOVE	044340
C		ARG=DFLOAT(INDEX-1)*PHI	044350
C		ARGUMENT OF SIN FN	044360
C		CRS=DCPLX(DRRES(INDEX),DMRES(INDEX))	044370
C		CCOSTM(II)=DCOS(ARG)*CRS*4*PI	044380
C		CSINTM(II)=DSIN(ARG)*CRS*4*PI	044390
C		CALC COEFF CONSTRUCTION TERMS	044400
C	370	CONTINUE	044410
C		DO 380 JJ=1,NPREJ	044420
C		LOOP TO REPLACE COEFFS FOR PREASSI	044430
C		GNED J TERMS	044440
C		J=JPREAS(JJ)*NX	044450
C		INDEX OF COEF BEING REPLACED	044460
C		INDEX=2*JJ-NINDEX	044470
			044480

C		SERIES INDEX FOR REPLACING TERM	04490
	GO TO (371,372),NSCOS		04500
C		SELECT PROPER SERIES COEFF ACCORDI	04510
C		NG TO Y SYMMETRY CONDITION	04520
C		NSCOS=2 INDICATES COSINES IN X	04530
C		CURRENT EQ	04540
371	CHDM(I,JJ)=-PI/(2*CS)*(0.00,1.00)**INDX*(CSINTM(JJ)-CSINTM(JJ+1))		04550
	CHDM(NPTCHS+I,JJ)=-PI/(2*CS)*(0.00,1.00)**INDX*(CCOSTM(JJ)+		04560
	1CCOSTM(JJ+1))		04570
	GO TO 380		04580
372	CHDM(I,JJ)=-PI/(2*CS)*(0.00,1.00)**(INDX+1)*(CCOSTM(JJ+1)-		04590
	1CCOSTM(JJ))		04600
	CHDM(I+NPTCHS,JJ)=-PI/(2*CS)*(0.00,1.00)**(INDX+1)*(CSINTM(JJ)+		04610
	1CSINTM(JJ+1))		04620
380	CONTINUE		04630
	DO 390 IT=1,NPREI		04640
	J=(NY-1)*NX+IPREAS(IT)+NPTCHS		04650
C		LOOP TO REPLACE COL'S FOR PREASSI	04660
	JJ=IT+NPREJ		04670
C		GNED I TERMS	04680
	INDX=2*(IT+NPREJ)-NINDX		04690
	GO TO (381,382),NSCOS		04700
381	CHDM(I,JJ)=-PI/(2*CS)*(0.00,1.00)**INDX*(CSINTM(IT+NPREJ)-		04710
	1CSINTM(IT+NPREJ+1))		04720
	CHDM(I+NPTCHS,JJ)=-PI/(2*CS)*(0.00,1.00)**INDX*(CCOSTM(IT+NPREJ)+		04730
	1CCOSTM(IT+NPREJ+1))		04740
	GO TO 390		04750
382	CHDM(I,JJ)=-PI/(2*CS)*(0.00,1.00)**(INDX+1)*(CCOSTM(IT+NPREJ+1)-		04760
	1CCOSTM(IT+NPREJ))		04770
	CHDM(I+NPTCHS,JJ)=-PI/(2*CS)*(0.00,1.00)**(INDX+1)*		04780
	1(CSINTM(IT+NPREJ)+CSINTM(IT+NPREJ+1))		04790
390	CONTINUE		04800
400	CONTINUE		04810
C			04820
C		END OF MOMENT MATRIX CONSTRUCTION	04830
C			04840
405	CONTINUE		04850
	CALL SPRHOM(CMATX,NPTCHS,NPTCHS-NPREJ,NDIM1,NDIM1,		04860
1	CMATY,NPTCHS,NPTCHS-NPREI,NDIM1,NDIM1,		04870
2	CHDM,NDIM1,NDIM1,CMAT4,NDIM1,NDIM1,CDET)		04880
	FRAT=CDABS(CMATX(1,1))		04890
	CPLATE=CDET/FRAT		04900
	WRITE(6,20) CSUNDR,CPLATE		04910
20	FORMAT(' ',5X,2F12.4,5X,2F12.4)		04920
	RETURN		04930
	END		04940

	SUBROUTINE SPRHOM(CMAT1,NI1,NJ1,NDIM1I,NDIM1J,CMAT2,NI2,NJ2,NDIM2I	00010
	1,NDIM2J,CMAT3,NDIM3I,NDIM3J,CMAT4,NDIM4I,NDIM4J,CDET)	00020
	IMPLICIT COMPLEX*16(C),REAL*8(A,B,D-H,C-Z)	00030
C		00040
C	SUBROUTINE TO DIAGCNALIZE AND CALC DETERMINANT OF A SPARCELY-	00050
C	COUPLED MATRIX	00060
C	BY L W PEARSON 7/74	00070
C	REVISED 5/75	00075
C		00080
	DIMENSION CMAT1(NDIM1I,NDIM1J),CMAT2(NDIM2I,NDIM2J),CMAT3(NDIM3I,	00090
	NDIM3J),CMAT4(NDIM4I,NDIM4J)	00100
	NI3=NI1+NI2	00110
	NJ3=NI3-NJ2-NJ1	00120
	CALL ZEROZ(CMAT4,4*NDIM4I*NDIM4J)	00130
	CDET=1	00140
C	INITIALIZE PRODUCT ACCUMULATOR	00150
	NPR=3	
	NJ1M1=NJ1-1	00160
	NJ1L=NJ1	
	IF(NJ2+NJ3.GE.1) GO TO 95	
	NJ1L=NJ1L-1	
	NPR=1	
95	DO 155 M=1,NJ1L	00170
C	INDEX ACROSS COL	00180
	MP1=M+1	00190
	FMAX=CDABS(CMAT1(M,M))	00200
	K=M	00210
	IF(MP1.GT.NI1) GO TO 105	00220
	DO 100 I=MP1,NI1	00230
C	LGCF TC SEARCH FOR PIVOT IN MTH	00240
C	COL	00250
	FCK=CDABS(CMAT1(I,M))	00260
	IF(FCK.LE.FMAX) GO TO 100	00270
	K=I	00280
C	IF LARGER ELEMENT FOUND MARK ROW	00290
	FMAX=FCK	00300
C	USE NEW LARGE ELEMENT AS COMPARI-	00310
C	SON VALUE	00320
100	CONTINUE	00330
105	CSTOR=CMAT1(K,M)	00340
C	SAVE VAL OF PIVOT ELEMENT	00350
	CDET=CDET*CSTOR	00360
C	MULT PIVOT INTO PROD ACCUMULATOR	00370
	IF(K.EQ.M) GO TO 115	00380
C	IF PIVOT CN DIAG SKIP ROW EXCH	00390
	CDET=-CDET	00400
C	CHANGE SIGN BECAUSE OF ROW EXCH	00410
107	DO 110 J=M,NJ1	00420
C	LOOP TO EXCH DIAG AND PIVOT ROWS	00430
	CSTC=CMAT1(K,J)	00440
	CMAT1(K,J)=CMAT1(M,J)	00450
	CMAT1(M,J)=CSTC	00460
110	CONTINUE	00470
	IF(NJ3.LT.1) GO TO 115	00475
	DO 112 J=1,NJ3	00480
	CSTO=CMAT3(K,J)	00490
	CMAT3(K,J)=CMAT3(M,J)	00500
	CMAT3(M,J)=CSTO	00510
112	CONTINUE	00520
115	CONTINUE	00560
	IF(MP1.GT.NI1) GO TO 155	00570

C	DO 150 I=MP1,NI1		00580
	CFAC=CMAT1(I,M)/CSTOR	ELIMINATION LOOP FOR CMAT1	00590
C		ELIMINATION FACTOR	00600
	IF(MP1.GT.NJ1) GO TO 125		00610
	DO 120 J=MP1,NJ1		00620
C		LOOP ACROSS ROW IN CMAT3	00630
	CMAT1(I,J)=CMAT1(I,J)-CMAT1(M,J)*CFAC		00640
120	CONTINUE		00650
	IF(NJ3.LT.1) GO TO 150		00660
125	DO 130 J=1,NJ3		00665
C		LOOP ACROSS ROW IN CMAT3	00670
	CMAT3(I,J)=CMAT3(I,J)-CMAT3(M,J)*CFAC		00680
130	CONTINUE		00690
150	CONTINUE		00700
155	CONTINUE		00720
	NI4=NI1-NJ1		00730
	IF(NI4.LE.0) GO TO 290		00740
C			00750
C			00760
C	BEGIN ROUTINE TO CREATE/'DIAGONALIZE' CMAT4		00770
C			00780
	NPIV=NI4		00790
	IF(NI4.GT.NJ2) NPIV=NJ2		00800
	DO 250 M=1,NPIV		00810
C		INDEX ACROSS COL FOR CMAT4 DIAG	00820
	MP1=M+1		00830
	FMAX=CCABS(CMAT2(1,M))		00840
	K=1		00850
	IF(NI2.LT.2) GO TO 205		00860
	DO 200 I=2,NI2		00870
C		LOOP TO SEARCH FOR PIVOT IN MTH	00880
C		COL	00890
	FCK=CCABS(CMAT2(I,M))		00900
	IF(FCK.LE.FMAX) GO TO 200		00910
	K=I		00920
C		IF LARGER ELEMENT FOUND MARK ROW	00930
	FMAX=FCK		00940
C		USE NEW LARGE ELEMENT AS COMPARI-	00950
C		SON VALUE	00960
200	CONTINUE		00970
205	CSTOR=CMAT2(K,M)		00980
C		SAVE VAL OF PIVOT ELEMENT	00990
	CDET=CDET*CSTOR		01000
C		MULT PIVOT INTO PROD ACCUM	01010
	CDET=-CDET		01020
C		CHANGE SIGN OF DETERM BECAUSE OF	01030
C		EXCHANGE FROM CMAT2 TO CMAT4	01040
	DO 210 J=M,NJ2		01050
C		LOOP TO EXCHANGE DIAG AND PIVCT R	01060
C		ROWS	01070
	CSTO=CMAT4(M,J)		01080
	CMAT4(M,J)=CMAT2(K,J)		01090
	CMAT2(K,J)=CSTO		01100
210	CONTINUE		01110
	K3=K+NI1		01120
	M3=NJ1+M		01130
	IF(NJ3.LT.1) GO TO 213		01140
	DO 212 J=1,NJ3		01145
	CSTO=CMAT3(K3,J)		01150
	CMAT3(K3,J)=CMAT3(M3,J)		01160
			01170

	CMAT3(M3,J)=CSTO	01180
212	CONTINUE	01190
213	IF(NI2.LT.1) GC TC 290	01225
235	DO 250 I=1,NI2	01230
C		01240
C	LOOP TO CARRY ELIMINATION INTO CMAT2	01250
	I3=NI1+1	01260
	CFAC=CMAT2(I,M)/CSTO	01270
	IF(MP1.GT.NJ2) GO TC 242	01280
	DO 240 J=MP1,NJ2	01290
C		01300
	LOOP ACROSS ROW OF CMAT2	01310
	CMAT2(I,J)=CMAT2(I,J)-CMAT4(M,J)*CFAC	01320
240	CONTINUE	01325
	IF(NJ3.LT.1) GO TO 250	01330
242	DO 245 J=1,NJ3	01340
C		01350
	LOOP ACROSS ROW OF CMAT3	01360
	CMAT3(I3,J)=CMAT3(I3,J)-CMAT3(M3,J)*CFAC	01380
245	CONTINUE	01390
250	CONTINUE	01400
C		01410
C	END ROUTINE TO 'DIAGONALIZE' CMAT4	01420
C		01430
290	IF(NI4.GE.NJ2) GO TO 350	01440
C		01450
C	IF DIAGONAL DOES NOT PASS THRU SKIP DIAGONALIZATION FOR CMAT2	01460
C		01470
C	BEGIN ROUTINE TO 'DIAGONALIZE' CMAT2	01480
	NI4P1=NI4+1	01482
	NJ2L=NJ2	01484
	IF(NJ3.GE.1) GC TC 295	01486
	NJ2L=NJ2L-1	01488
	NPR=2	01492
295	DO 350 M=NI4P1,NJ2L	01500
	MI=M-NI4	01510
	M3=MI+NI1	01515
	MP1=M+1	01520
	MIP1=MI+1	01530
	FMAX=CDABS(CMAT2(MI,M))	01540
	K=MI	01550
	IF(MIP1.GT.NI2) GO TO 305	01560
	DO 300 I=MIP1,NI2	01570
C		01580
C	LOOP TO SEARCH FOR PIVOT IN MTH COL	01590
	FCK=CDABS(CMAT2(I,M))	01600
	IF(FCK.LE.FMAX) GO TO 300	01610
	K=I	01620
C		01630
	IF LARGER ELEMENT FOUND MARK ROW	01640
	FMAX=FCK	01650
C		01660
C	USE NEW LARGE ELEMENT AS COMPARI- SON VALUE	01670
300	CONTINUE	01680
305	CSTO=CMAT2(K,M)	01690
C		01700
	SAVE VAL OF PIVOT ELEMENT	01710
	K3=K+NI1	01720
	CDET=CDET*CSTO	01730
C		01740
	MULT PIVOT INTO PROD ACCUMULATOR	01750
	IF(K.EQ.MI) GO TC 315	
C		
	IF PIVOT ON DIAG SKIP ROW EXCH	
	CDET=-CDET	
C		
	CHANGE SIGN BECAUSE OF ROW EXCH	

	DO 310 J=M,NJ2		C1760
	CSTO=CMAT2(K,J)		01770
	CMAT2(K,J)=CMAT2(MI,J)		01780
	CMAT2(MI,J)=CSTO		01790
310	CONTINUE		01800
	IF(NJ3.LT.1) GO TC 315		C1805
	DO 312 J=1,NJ3		01810
C		LOOP FOR EXCH IN CMAT3	01820
	CSTO=CMAT3(K3,J)		01830
	CMAT3(K3,J)=CMAT3(M3,J)		01840
	CMAT3(M3,J)=CSTO		C1850
312	CONTINUE		01860
315	CONTINUE		01900
	IF(MIPL.GT.NI2) GO TO 390		C1910
	DO 350 I=MIPL,NI2		01920
C		ELIMINATION LOOP	01930
	I3=I+NI1		C1940
	CFAC=CMAT2(I,M)/CSTO		01950
	IF(MIPL.GT.NJ2) GO TO 335		01960
	DO 330 J=MIPL,NJ2		C1970
C		LOOP ACROSS ROW OF CMAT2	01980
	CMAT2(I,J)=CMAT2(I,J)-CMAT2(MI,J)*CFAC		01990
330	CONTINUE		02000
335	IF(NJ3.LT.1) GO TC 350		02005
	DO 345 J=1,NJ3		02010
C		LOOP ACROSS ROW IN CMAT3	02020
	CMAT3(I3,J)=CMAT3(I3,J)-CMAT3(M3,J)*CFAC		02030
345	CONTINUE		02040
350	CONTINUE		C2060
C			02070
C	BEGIN ROUTINE TO 'DIAGONALIZE' CMAT3		C2080
C			02090
390	NJ3M1=NJ3-1		02100
	IF(NJ3M1.LT.1) GO TC 455		02110
	DO 450 M=1,NJ3M1		02120
C		INDEX ACROSS COL	02130
	MIPL=M+1		02140
	MI=M+NJ1+NJ2		02150
	MIPL=MI+1		C2160
	FMAX=CDABS(CMAT3(MI,M))		02170
	K=MI		02180
	IF(MIPL.GT.NI3) GO TO 405		C2190
	DO 400 I=MIPL,NI3		02200
C		LOOP TO SEARCH FOR PIVCT IN MTH	02210
C		COL	C2220
	FCK=CDABS(CMAT3(I,M))		02230
	IF(FCK.LE.FMAX) GO TO 400		02240
	K=I		02250
C		IF LARGER ELEMENT FOUND MARK ROW	02260
	FMAX=FCK		C2270
C		USE NEW LARGE ELEMENT AS COMPARI-	02280
C		SO VALUE	02290
400	CONTINUE		C2300
405	CSTO=CMAT3(K,M)		02310
C		SAVE VAL OF PIVOT ELEMENT	02320
	CDET=CDET*CSTO		C2330
C		MULT PIVOT INTO PROD ACCUMULATOR	02340
	IF(K.EQ.MI) GO TO 415		C2350
C		IF PIVCT ON DIAG SKIP ROW EXCH	02360
	CDET=-CDET		02370
C		CHANGE SIGN BECAUSE OF ROW EXCH	C2380

C	DO 410 J=M,NJ3		02390
		LOOP TO EXCH DIAG AND PIVCT RCWS	02400
	CSTO=CMAT3(K,J)		02410
	CMAT3(K,J)=CMAT3(MI,J)		02420
	CMAT3(MI,J)=CSTO		02430
410	CONTINUE		02440
415	CONTINUE		02480
	DO 450 I=MIP1,NI3		02490
C		ELIMINATION LCOP	02500
	CFAC=CMAT3(I,M)/CSTOR		02510
	DO 445 J=MPI,NJ3		02520
C		LOOP ACROSS ROW IN CMAT3	02530
	CMAT3(I,J)=CMAT3(I,J)-CMAT3(MI,J)*CFAC		02540
445	CONTINUE		02550
450	CONTINUE		02570
455	GO TO (461,462,463), NFR		02572
461	CDET=CDET*CMAT1(NI1,NJ1)		02574
	RETURN		02576
462	CDET=CDET*CMAT2(NI2,NJ2)		02578
	RETURN		02582
463	CDET=CDET*CMAT3(NI3,NJ3)		02584
	RETURN		02600
C		MULT LAST ELEMENT INTC DETERM	02590
	END		02610

	SUBROUTINE SPRSLV (CMAT1,NI1,NJ1,NDIM1I,NDIM1J,CMAT2,NI2,NJ2,NDIM2I	00010
	1,NDIM2J,CMAT3,NDIM3I,NDIM3J,CMAT4,NDIM4I,NDIM4J,CRHS,C SOLN)	00020
C		00030
C	SUBROUTINE TO BACKSLVE A TRIANGULARIZED SYSTEM OF SPARCELY-	00040
C	COUPLED LINEAR EQUATION	00050
C	BY L W PEARSON 7/74	00052
C	REVISED 5/75	00054
C		00056
C	STORAGE FORM COMPATIBLE WITH THE TRIANGULARIZATION ROUTINE SPARCE	00060
C		00070
C	THE ENTRY 'HOMSLV' BELOW ALLOWS THE SOLUTION FOR NATURAL VECTORS	00080
C	OF HOMOGENEOUS SYSTEMS PROVIDED THE DETERMINANT OF THE SYSTEM IS	00090
C	ZERO	00100
C		00110
	IMPLICIT COMPLEX*16(C),REAL*8(A,B,D-H,C-Z)	00120
	DIMENSION CMAT1(NDIM1I,NDIM1J),CMAT2(NDIM2I,NDIM2J),CMAT3(NDIM3I,N	00130
	NDIM3J),CMAT4(NDIM4I,NDIM4J),CRHS(NDIM3I),C SOLN(NDIM3I)	00140
	LOGICAL LHOM	00150
C		00160
C	SETUP FOR INHOMOGENEOUS SYSTEM	00170
C		00180
	LHOM=.FALSE.	00190
C	SET INDICATOR FOR INHOM ENTRY	00200
	NI3=NI1+NI2	00210
C	NO ROWS IN COUPLING SUBMATRIX	00220
	NJ3=NJ1-NJ2	00230
C	NO CF COLS	00240
	NI4=NI1-NJ1	00250
C	NO ROWS IN SECONDARY COUPLING	00260
C	SUBMATRIX	00270
	ND2=NJ2-NI4	00280
C	NO OF DIAGONAL TERMS OF MATRIX	00290
C	IN CMAT2	00300
	NPR=3	
	IF(NJ3.LT.1) NPR=2	
C	SET INDICATOR FOR NULL CMAT3	
C	DEGENERACY	
	IF(NJ3+NJ2.LT.1) NPR=1	
C	SET INDICATOR FOR NULL CMAT2 &	
C	CMAT3	
	GO TO (81,82,83) , NPR	
C	GO MAKE FIRST DIVISION FOR RIGHT-	
C	MOST MATRIX	
81	C SOLN(NI3)=CRHS(NI3)/CMAT1(NI1,NJ1)	
	GO TO 100	
82	C SOLN(NI3)=CRHS(NI3)/CMAT2(NI2,NJ2)	
	GO TO 100	
83	C SOLN(NI3)=CRHS(NI3)/CMAT3(NI3,NJ3)	
C	SOLVE FOR 'LAST' UNKNOWN	00320
	GO TO 100	00330
C	GO TO SOLN ROUTINES	00340
C		00350
C	END OF SETUP FOR INHOM SYSTEM	00360
C		00370
C	BEGIN ENTRY/SETUP FOR HOMOGENEOUS SYSTEM	00380
C		00390
	ENTRY HOMSLV(CMAT1,NI1,NJ1,NDIM1I,NDIM1J,CMAT2,NI2,NJ2,NDIM2I,NDIM	00400
	12J,CMAT3,NDIM3I,NDIM3J,CMAT4,NDIM4I,NDIM4J,C SOLN,NORD)	00410
C		00420
	LHOM=.TRUE.	00430
C	LOGICAL INDICATOR FOR HOMOGEN SYS	00440

	NI3=NI1+NI2	00450
	NJ3=NI3-NJ1-NJ2	00460
	NI4=NJ1-NJ1	00470
	ND2=NJ2-NI4	00480
	CSOLN(NI3)=1	00490
C	ASSIGN ARBITRARY ELEMENT IN SCL'N	00500
C		00510
C	END SETUP FOR HOMOGENEOUS ENTRY	00520
C		00530
C	BEGIN BACKSOLVE FOR EQUATIONS INVOLVING ONLY CMAT3 (LAST NJ3 EQS)	00540
C		00550
100	FMAX=CDABS(CSOLN(NI3))	00560
	IMAX=NI3	00570
	IF(NJ3.LT.2) GO TO 200	00580
C	SKIP ROUTINE IF ONLY LAST VARIABLE	00590
C	COUPLES (IT WAS SOLVED/ASSIGNED	00600
C	ABOVE)	00610
		00620
	DO 150 IC=2,NJ3	00630
	ICM1=IC-1	00640
	I=NI3-IC+1	00650
	I=NI3-IC+1	00660
C	CALC MATRIX ROW INDX FROM	00670
C	COMPLEMENTARY INDX	00680
		00690
C	JD3=I-NJ1-NJ2	00700
C		00710
	COL INDX FOR CMAT3 WHICH DEFINES	00720
	DIAG OF MATRIX	00730
		00740
	CSUM=0	00750
	DO 110 J3C=1,ICM1	00760
C	LOOP TO ACCUM NEGATIVE SUM OF	00770
C	PREVIOUSLY CALC'D UNKNS	00780
		00790
C	J3=NJ3+1-J3C	00800
C		00810
	COL OF COEF IN CMAT3	00820
C	J=NI3+1-J3C	00830
C		00840
	ROW OF UNKN IN CSCLN	00850
		00860
	CSUM=CSUM-CMAT3(I,J3)*CSOLN(J)	00870
110	CONTINUE	00880
	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I)	00890
C	ADD R H S TO SUM	00900
		00910
C	CSOLN(I)=CSUM/CMAT3(I,JD3)	00920
C	DIVIDE BY DIAG COEF	00930
		00940
	IF(CDABS(CSOLN(I)).LE.FMAX) GO TO 150	00950
	FMAX=CDABS(CSCLN(I))	00960
	IMAX=I	00970
C	CHECK FOR MAX ELEMENT	00980
150	CONTINUE	00990
C		01000
C	BEGIN ROUTINE TO SOLVE FOR ELEMENTS INVOLVING CMAT3 & CMAT2	01010
C		01020
200	IF(NJ3.GE.NI2) GO TO 300	01030
C	SKIP ROUTINE IF DIAG DOES NOT	01040
C	PASS THRU CMAT2	
	DO 250 IC=1,ND2	
	ICM1=IC-1	
	I2=NI2-NJ3+1-IC	
	I3=NI3-NJ3+1-IC	
	JD2=NJ2+1-IC	
	NCM1=NJ3+IC-1	
	CSUM=0	
	IF(NJ3.LT.1) GO TO 215	
	DO 210 JC=1,NJ3	
C	LOOP TO SUM CONTRIB FROM CMAT3	

	J3=NJ3+1-JC	01050
	J=NI3+1-JC	01060
	CSUM=CSUM-CMAT3(I3,J3)*CSOLN(J)	01070
210	CONTINUE	01080
215	IF(ICM1.LT.1) GO TO 225	01090
C	SKIP IF NO TERMS CONTRIB FR CMAT2	01100
	DO 220 J2C=1,ICM1	01110
	J2=NJ2+1-J2C	01120
	J=NI3-NJ3+1-J2C	01130
	CSUM=CSUM-CMAT2(I2,J2)*CSOLN(J)	01140
220	CONTINUE	01150
225	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I3)	01160
	CSOLN(I3)=CSUM/CMAT2(I2,J2)	01170
	IF(CDABS(CSOLN(I3)).LE.FMAX) GO TO 250	01180
	FMAX=CDABS(CSCLN(I3))	01190
	IMAX=I	01200
250	CONTINUE	01210
C		01220
C	BEGIN ROUTINE TO SOLVE FOR ELEMENTS INVOLVING CMAT3 & CMAT4	01230
C		01240
300	IF(NI4.LT.1) GO TO 400	
	DO 350 IC=1,NI4	
	I4=NI4+1-IC	01260
	JD4=I4	01270
	I3=NI1+1-IC	01280
	CSUM=0	01290
	IF(NJ3.LT.1) GO TO 315	
	DO 310 J3C=1,NJ3	01300
	J3=NJ3+1-J3C	01310
	J=NI3+1-J3C	01320
	CSUM=CSUM-CMAT3(I3,J3)*CSOLN(J)	01330
310	CONTINUE	01340
315	NSUBS=ND2+IC-1	01350
C	NO CF NON-DIAG CMAT4 EL'S IN EQ	01360
	IF(NSUBS.LT.1) GO TO 325	01370
	DO 320 J4C=1,NSUBS	01380
	J4=NJ2+1-J4C	01390
	J=NI3-NJ3+1-J4C	01400
	CSUM=CSUM-CMAT4(I4,J4)*CSOLN(J)	01410
320	CONTINUE	01420
325	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I3)	01430
	CSOLN(I3)=CSUM/CMAT4(I4,I4)	01440
	IF(CDABS(CSOLN(I3)).LE.FMAX) GO TO 350	01450
	FMAX=CDABS(CSCLN(I3))	01460
	IMAX=I3	01470
350	CONTINUE	01480
C		01490
C	BEGIN ROUTINE TO SOLVE EQ'S INVOLVING CMAT3 & CMAT1	01500
C		01510
400	IF(NJ1.LT.1) GO TO 455	
	DO 450 IC=1,NJ1	01520
	I=NJ1+1-IC	01530
	ICM1=IC-1	01540
	CSUM=0	01550
	IF(NJ3.LT.1) GO TO 415	
	DO 410 J3C=1,NJ3	01560
	J3=NJ3+1-J3C	01570
	J=NI3+1-J3C	01580
	CSUM=CSUM-CMAT3(I,J3)*CSOLN(J)	01590
410	CONTINUE	01600
415	IF(ICM1.LT.1) GO TO 425	01610

	DO 420 JC=1, ICM1	01620
	J=NJ1+1-JC	01630
	CSUM=CSUM-CMAT1(I,J)*CSOLN(J)	01640
420	CONTINUE	01650
425	IF(.NOT.LHOM) CSUM=CSUM+CRHS(I)	01660
	CSOLN(I)=CSUM/CMAT1(I,I)	01670
	IF(CDABS(CSOLN(I)).LE.FMAX) GO TO 450	01680
	FMAX=CDABS(CSOLN(I))	01690
	IMAX=I	01700
450	CONTINUE	01710
C		01720
C		01730
C	END OF SOLUTION	01740
C		01750
455	IF(.NOT.LHOM) RETURN	01760
C		01770
C	RETURN IF INHOM SYSTEM	01780
C		01790
C	BEGIN NORMALIZATION ROUTINE FOR NATURAL VECTOR FOR HOMOGENEOUS	01800
C	CASE	01810
	CSCALE=1./CSOLN(IMAX)	01820
	DO 500 I=1, NI3	01830
	CSOLN(I)=CSOLN(I)*CSCALE	01840
500	CONTINUE	01850
	RETURN	01860
	END	01870

SUBROUTINE COPYZ(X,Y,N)	09340
DIMENSION X(1),Y(1)	09350
DO 100 I=1,N	09360
X(I)=Y(I)	09370
100 CONTINUE	09380
RETURN	09390
END	09400

```
100 SUBROUTINE ZEROZ(IARRAY,N)
    DIMENSION IARRAY(1)
    DO 100 I=1,N
    IARRAY(I)=0
    CONTINUE
    RETURN
    END
```

```
09410
09420
09430
09440
09450
09460
09470
```

	SUBROUTINE DWEDDL(FCN,N,DELTA,VINT)	09480
	IMPLICIT REAL*8(A-H,O-Z)	09490
	COMPLEX*16 FCN,C,VINT	09500
	DIMENSION FCN(N)	09510
	DIMENSION COEF(6)	09520
	DATA COEF/2.00,5.00,1.00,6.00,1.00,5.00/	09530
	IF((N-1)/6*6.EQ.N-1) GO TO 100	09540
	WRITE(6,1)	09550
1	FORMAT('OINCORRECT POINTS TO WEDDLE')	09560
	A=1/O	09570
100	CONTINUE	09580
	VINT=0	09590
	DO 200 J=1,N	09600
	JCOEF=J-((J-1)/6)*6	09610
	VINT=VINT+COEF(JCOEF)*FCN(J)	09620
200	CONTINUE	09630
	VINT=(VINT-FCN(1)-FCN(N))*(0.300,0.00)*DCMPLX(DELTA,0.00)	09640
	RETURN	09650
	END	09660

C	ZANLYT.....	ZAN09670
C			ZAN09680
C	FUNCTION	- DETERMINATION OF ZEROS OF AN ANALYTIC COMPLEX FUNCTION USING MULLER'S METHOD WITH DEFLATION	ZAN09690 ZAN09700 ZAN09710
C	USAGE	- CALL ZANLYT (F,EPS,NSIG,KN,NGUESS,N,X,ITMAX,INFER,IER)	ZAN09720 ZAN09730
C	PARAMETERS	F	ZAN09740
C		- A FUNCTION SUBPROGRAM, F(Z), WRITTEN BY THE USER SPECIFYING THE EQUATION WHOSE ROOTS ARE TO BE FOUND. F MUST BE TYPE-NAMED AS FOLLOWS - COMPLEX FUNCTION F*16 (Z)	ZAN09750 ZAN09760 ZAN09770
C		EPS	ZAN09780
C		- 1ST STOPPING CRITERION. A ROOT Z IS ACCEPTED IF ABSOLUTE VALUE OF F(Z) .LE. EPS (INPUT)	ZAN09790
C		NSIG	ZAN09800
C		- 2ND STOPPING CRITERION. A ROOT IS ACCEPTED IF TWO SUCCESSIVE APPROXIMATIONS TO A GIVEN ROOT AGREE IN THE FIRST NSIG DIGITS. (INPUT)	ZAN09810 ZAN09820
C			ZAN09830
C		NOTE. IF EITHER OR BOTH OF THE STOPPING CRITERIA ARE FULFILLED, THE ROOT IS ACCEPTED.	ZAN09840 ZAN09850
C		KN	ZAN09860
C		- THE NUMBER OF KNOWN ROOTS WHICH MUST BE STORED IN X(1),...,X(KN), PRIOR TO ENTRY TO ZANLYT	ZAN09870
C		NGUESS	ZAN09880
C		- THE NUMBER OF INITIAL GUESSES PROVIDED. THESE GUESSES MUST BE STORED IN X(KN+1),...,X(KN+NGUESS) AND NGUESS MUST BE SET EQUAL TO ZERO IF NO GUESSES ARE PROVIDED. (INPUT)	ZAN09890 ZAN09900 ZAN09910
C		N	ZAN09920
C		- THE NUMBER OF NEW ROOTS TO BE FOUND BY ZANLYT (INPUT)	ZAN09930
C		X	ZAN09940
C		- A LONG-WORD COMPLEX VECTOR ARRAY OF LENGTH .GE. 3*(KN+N). X(1),...,X(KN) ON INPUT MUST CONTAIN ANY KNOWN ROOTS. X(KN+1),...,X(KN+N) ON INPUT MAY, AT THE USER'S OPTION, CONTAIN INITIAL GUESSES FOR THE N NEW ROOTS WHICH ARE TO BE COMPUTED. ON OUTPUT, X(KN+1),...,X(KN+N) CONTAIN EITHER A ROOT CORRECT TO WITHIN A CONVERGENCE CRITERION OR THE VALUE(12345678.12345678D+0,12345678.12345678D+0) INDICATIVE OF A FAILURE TO ACHIEVE THE SPECIFIED CONVERGENCE FOR THAT ROOT, SAY X(KN+J). IN THE LATTER CASE, THE MOST RECENT APPROXIMATION TO X(KN+J) IS AVAILABLE IN X(ISUB), WHERE ISUB=2*(KN+N)+J	ZAN09950 ZAN09960 ZAN09970 ZAN09980 ZAN09990 ZAN10000 ZAN10010 ZAN10020 ZAN10030 ZAN10040 ZAN10050 ZAN10060 ZAN10070
C		ITMAX	ZAN10080
C		- THE MAXIMUM ALLOWABLE NUMBER OF ITERATIONS PER ROOT (INPUT)	ZAN10090
C		INFER	ZAN10100
C		- AN INTEGER VECTOR OF LENGTH .GE. KN+N. ON OUTPUT INFER(J) CONTAINS THE NUMBER OF ITERATIONS USED IN FINDING THE J-TH ROOT WHEN CONVERGENCE WAS ACHIEVED. IF CONVERGENCE WAS NOT OBTAINED IN ITMAX ITERATIONS, INFER(J) WILL CONTAIN ITMAX+1 (OUTPUT)	ZAN10110 ZAN10120 ZAN10130 ZAN10140 ZAN10150
C		IER	ZAN10160
C		- ERROR PARAMETER (OUTPUT)	ZAN10170
C		WARNING ERROR = 32 + N	ZAN10180
C		N = 1 FAILURE TO CONVERGE WITHIN ITMAX ITERATIONS FOR ONE OF THE (N) NEW ROOTS TO BE FOUND	ZAN10190 ZAN10200
C			ZAN10210
C	PRECISION	- DOUBLE	ZAN10220
C	REQ'D IMSL ROUTINES	- UERTST	ZAN10230
C	AUTHOR/IMPLEMENTOR	- C. G. JOHNSON/L. L. WILLIAMS	ZAN10240
C	LANGUAGE	- FORTRAN	ZAN10250
C			ZAN10260
C	LATEST REVISION	- SEPTEMBER 1, 1971	ZAN10270

C	SUBROUTINE ZANLYT (F, FPS, NSIG, KN, NGUESS, N, X, ITMAX, INFER, IER)	ZAN10280
	COMPLEX*16 X(1), ONE, D, DD, DEN, DI, FPRT, FRT,	ZAN10290
1	H, RT, T1, T2, T3, TEM, XO, X1, X2, BI, F, XX	ZAN10300
	DOUBLE PRECISION QZ, FPS, FPS1	ZAN10310
	DIMENSION INFER(1)	ZAN10320
	IER = 0	ZAN10330
	ONE = (1.0D+00, 0.0D+00)	ZAN10340
	FPS1 = 10.0D+00**(-NSIG)	ZAN10350
	ICONJ = 0	ZAN10360
	IRDMB = 0	ZAN10370
C	SET NUMBER OF ITERATIONS	ZAN10380
	MB1 = KN+1	ZAN10390
	MB2 = KN+N	ZAN10400
	LSTART = MB2+1	ZAN10410
	MPG = MB1+NGUESS	ZAN10420
	DO 2 I = MPG, MB2	ZAN10430
2	X(I) = (0.0D+0, 0.0D+0)	ZAN10440
	L = MB1	ZAN10450
	IF (KN .EQ. 0) GO TO 5	ZAN10460
	DO 3 I = 1, KN	ZAN10470
	INFER(I) = 0	ZAN10480
	ITEMP = MB2+I	ZAN10490
	X(ITEMP) = X(I)	ZAN10500
	ITEMP = MB2+ITEMP	ZAN10510
3	X(ITEMP) = X(I)	ZAN10520
5	JK = 0	ZAN10530
	QZ = CDABS(X(L))	ZAN10540
	IF (QZ .LE. 1.0D-15) GO TO 25	ZAN10550
C	ROOT ESTIMATE NOT EQUAL TO ZERO	ZAN10560
10	RT = (.9D+00, 0.0D+00)*X(L)	ZAN10570
	ASSIGN 15 TO NN	ZAN10580
	GO TO 135	ZAN10590
15	XO = FPRT	ZAN10600
	RT = (1.1D+00, 0.0D+00)*X(L)	ZAN10610
	ASSIGN 20 TO NN	ZAN10620
	GO TO 135	ZAN10630
20	X1 = FPRT	ZAN10640
	H = X(L)-RT	ZAN10650
	RT = X(L)	ZAN10660
	ASSIGN 40 TO NN	ZAN10670
	GO TO 135	ZAN10680
C	ROOT ESTIMATE EQUAL TO ZERO	ZAN10690
25	RT = -ONE	ZAN10700
	ASSIGN 30 TO NN	ZAN10710
	GO TO 135	ZAN10720
30	XO = FPRT	ZAN10730
	RT = ONE	ZAN10740
	ASSIGN 35 TO NN	ZAN10750
	GO TO 135	ZAN10760
35	X1 = FPRT	ZAN10770
	RT = (0.0D+00, 0.0D+00)	ZAN10780
	H = -ONE	ZAN10790
	ASSIGN 40 TO NN	ZAN10800
	GO TO 135	ZAN10810
40	X2 = FPRT	ZAN10820
45	D = (-0.5D+00, 0.0D+00)	ZAN10830
C	BEGIN MAIN ALGORITHM	ZAN10840
50	DD = ONE + D	ZAN10850
	T1 = XO*DD*D	ZAN10860
	T2 = X1*DD*DD	ZAN10870
		ZAN10880

XX = X2*DD	ZAN10890
T3 = X2*D	ZAN10900
R1 = T1-T2+XX+T3	ZAN10910
DEN = R1*R1-(4.00+0.0+0.0)*(XX*T1-T3*(T2-XX))	ZAN10920
USE DENOMINATOR OF MAXIMUM AMPLITUDE	ZAN10930
T1 = CDSQRT(DEN)	ZAN10940
T2 = R1 + T1	ZAN10950
T3 = R1 - T1	ZAN10960
OZ = CDABS(T2) - CDABS(T3)	ZAN10970
IF (OZ .GE. 1) GO TO 60	ZAN10980
55 DEN = T3	ZAN10990
GO TO 65	ZAN11000
60 DEN = T2	ZAN11010
TEST FOR ZERO DENOMINATOR	ZAN11020
65 OZ = CDABS(DEN)	ZAN11030
IF (OZ .GT. 1.0-15) GO TO 75	ZAN11040
70 DEN = ONE	ZAN11050
75 DT = ((-2.00+0.0+0.0)*XX)/DEN	ZAN11060
H = DT * H	ZAN11070
RT = RT + H	ZAN11080
CHECK CONVERGENCE OF THE FIRST KIND	ZAN11090
OZ = CDABS(H/PT)	ZAN11100
IF (OZ .LE. EPS1) GO TO 100	ZAN11110
80 ASSIGN 85 TO MN	ZAN11120
GO TO 135	ZAN11130
85 OZ = CDABS(FPRT)-CDABS(X2*(1.000+0.000))	ZAN11140
IF (OZ .LT. 0.0+0) GO TO 95	ZAN11150
TAKE REMEDIAL ACTION TO INDUCE CONVERGENCE	ZAN11160
90 DT = DT*(0.50+0.0+0.0)	ZAN11170
H = H*(0.50+0.0+0.0)	ZAN11180
RT = RT-H	ZAN11190
GO TO 135	ZAN11200
95 X0 = X1	ZAN11210
X1 = X2	ZAN11220
X2 = FPRT	ZAN11230
D = DT	ZAN11240
GO TO 50	ZAN11250
A ROOT HAS BEEN FOUND	ZAN11260
100 FRT = F(RT)	ZAN11270
105 X(L) = RT	ZAN11280
ITEMP = MR2+L-[ICMP	ZAN11290
X(ITEMP) = RT	ZAN11300
ITEMP = MR2+MR2+L	ZAN11310
X(ITEMP) = RT	ZAN11320
CHECK TO SEE IF COMPLEX-CONJUGATE IS ALSO A ROOT	ZAN11330
IF (CDABS(F(DCONJG(X(L)))) .GT. 10.0+0*CDABS(FRT)) GO TO 115	ZAN11340
OZ = CDABS(X(L)-DCONJG(X(L)))	ZAN11350
IF (ICONJ .NE. 0 .OR. OZ .LT. 1.00-8) GO TO 115	ZAN11360
ISTART = L+2	ZAN11370
INSERT1 = L+1	ZAN11380
DO 110 INSERT = ISTART,MR2	ZAN11390
X(INSERT) = X(INSERT1)	ZAN11400
110 INSERT = INSERT	ZAN11410
X(L+1) = DCONJG(X(L))	ZAN11420
ICONJ = 1	ZAN11430
GO TO 120	ZAN11440
115 ICONJ = 0	ZAN11450
120 CONTINUE	ZAN11460
125 INFER(L) = JK	ZAN11470
	ZAN11480
	ZAN11490

L = L+1		ZAN11500
IF (L .LE. MR2) GO TO 5		ZAN11510
	RETURN TO CALLING PROGRAM	ZAN11520
130 GO TO 185		ZAN11530
135 JK = JK+1		ZAN11540
IF (JK .GT. ITMAX) GO TO 180		ZAN11550
140 FRT = F(RT)		ZAN11560
FPRT = FRT		ZAN11570
	TEST TO SEE IF FIRST ROOT IS BEING DETERMINED	ZAN11580
		ZAN11590
IF (L .EQ. 1) GO TO 160		ZAN11600
IF (L .LE. IROMR+1) GO TO 160		ZAN11610
	COMPUTE DENOMINATOR FOR MODIFIED FUNCTION	ZAN11620
145 LIMUP = MR2+L-IROMR-1		ZAN11630
DO 150 I = LSTART,LIMUP		ZAN11640
TEM = RT - X(I)		ZAN11650
QZ = CDABS(TEM)		ZAN11660
IF (QZ .LT. 5.00-15) GO TO 175		ZAN11670
150 FPRT = FPRT/TEM		ZAN11680
	CHECK CONVERGENCE OF THE SECOND KIND	ZAN11690
160 QZ = CDABS(FRT)		ZAN11700
IF (QZ .GE. EPS) GO TO 170		ZAN11710
165 QZ = CDABS(FPRT)		ZAN11720
IF (QZ .LT. EPS) GO TO 105		ZAN11730
170 GO TO NN,(15,20,30,35,40,85)		ZAN11740
175 RT = RT * (1.0000010+0,0.00+0)		ZAN11750
GO TO 135		ZAN11760
	WARNING ERROR, ITMAX = MAXIMUM	ZAN11770
180 IFR = 33		ZAN11780
INFER(L) = ITMAX + 1		ZAN11790
IROMR = IROMR + 1		ZAN11800
X(L) = (12345678.123456780+0,12345678.123456780+0)		ZAN11810
ITEMP = MR2 + MR2 + L		ZAN11820
X(ITEMP) = RT		ZAN11830
L = L+1		ZAN11840
IF (L .LE. MR2) GO TO 5		ZAN11850
185 IF (IFR .EQ. 0) GO TO 9005		ZAN11860
9000 CONTINUE		ZAN11870
CALL UERTST(IFR,'ZANLYT')		ZAN11880
9005 RETURN		ZAN11890
END		ZAN11900
		ZAN11910


```

C.UERTST.....UEP11920
C      FUNCTION      - ERROR MESSAGE GENERATION      UEP11930
C      USAGE         - CALL UERTST(IEP,'NAMEXX')      UEP11940
C      PARAMETERS    IEP - ERROR PARAMETER. TYPE + N WHERE UEP11950
C                      TYPE= 128 IMPLIES TERMINAL ERROR UEP11960
C                      64 IMPLIES WARNING WITH FIX    UEP11970
C                      32 IMPLIES WARNING             UEP11980
C                      N = ERROR CODE RELEVANT TO CALLING ROUTINE UEP11990
C                      NAMEXX - NAME OF THE CALLING ROUTINE UEP12000
C      AUTHOR/IMPLEMENTER - PERER SVENDSEN          UEP12010
C      LANGUAGE       - FORTRAN                     UEP12020
C.....UEP12030
C      LATEST REVISION - JANUARY 19, 1971          UEP12040
C                      UEP12050
C      SUBROUTINE UERTST(IEP,NAME)                UEP12060
C      DIMENSION      ITYP(5,4),IRIT(4)          UEP12070
C      INTEGER*2      NAME(3)                    UEP12080
C      INTEGER        WARN,WARNF,TERM,PRINTR      UEP12090
C      EQUIVALENCE    (IRIT(1),WARN),(IRIT(2),WARNF),(IRIT(3),TERM) UEP12100
C      DATA          ITYP /'WARN','ING',' ',' ',' ',' ',' ' UEP12110
C      *              'WARN','ING','WITH','FIX',' ',' ' UEP12120
C      *              'TERM','INAL',' ',' ',' ',' ',' ' UEP12130
C      *              'NON-','DEFIN','ED',' ',' ',' ',' ' UEP12140
C      *              IRIT / 32,64,128,0/          UEP12150
C      DATA          PRINTR / 6/                 UEP12160
C      IF2=IEP                                       UEP12170
C      IF (IF2 .GE. WARN) GO TO 5                   UEP12180
C                      NON-DEFINED                UEP12190
C      IER1=4                                       UEP12200
C      GO TO 20                                     UEP12210
C      5 IF (IER2 .LT. TERM) GO TO 10              UEP12220
C                      TERMINAL                   UEP12230
C      IER1=3                                       UEP12240
C      GO TO 20                                     UEP12250
C      10 IF (IER2 .LT. WARNF) GO TO 15            UEP12260
C                      WARNING(WITH FIX)          UEP12270
C      IER1=2                                       UEP12280
C      GO TO 20                                     UEP12290
C                      WARNING                    UEP12300
C      15 IER1=1                                     UEP12310
C                      EXTRACT INI                UEP12320
C      20 IER2=IER2-IRIT(IER1)                    UEP12330
C                      PRINT ERROR MESSAGE        UEP12340
C      WRITE (PRINTR,25) (ITYP(I,IER1),I=1,5),NAME,IER2 UEP12350
C      25 FORMAT(' *** I M S L(UERTST) *** ',5I4,4X,3A2,4X,I2) UEP12360
C      RETURN                                       UEP12370
C      END                                          UEP12380
C                                          UEP12390
C                                          UEP12400

```

Reproduced from
best available copy.

	SUBROUTINE RSLJZ(X,FJ,NMAX,A,ND,TEPS,FJAPRX,RE)	00012410
	IMPLICIT REAL*8 (A-H,O-Z)	00012420
	DIMENSION FJ(1),FJAPRX(1),EP(1)	00012430
	NMAXT=NMAX	00012440
	IF(NMAXT.GE.0)GO TO 30	00012450
	IF(DABS(A).LE.1.0D-15)GO TO 10	00012460
	GO TO 20	00012470
10	IERR=4	00012480
	RETURN	00012490
20	NMAX0R=IABS(NMAXT)	00012500
	NMAXT=1	00012510
30	IF(A.GT.0.0)GO TO 40	00012520
	IF(DABS(A).LE.1.0D-15)GO TO 40	00012530
	IERR=1	00012540
	RETURN	00012550
40	IF(A.LT.1.0)GO TO 70	00012560
	IERR=2	00012570
	RETURN	00012580
70	IF(X.GT.0.0)GO TO 130	00012590
	IERR=3	00012600
	RETURN	00012610
130	TEPS=0	00012620
	EPSILON=.500*10.**(-ND)	00012630
	NMP1=NMAX+1	00012640
	DO 160 N=1,NMP1	00012650
160	FJAPRX(N)=0.0	00012660
	SJM=(X/2.):**A/DGAMMA(1.+A)	00012670
	D1=2.302600*ND+1.386300	00012680
	IF(NMAXT.LE.0)GO TO 230	00012690
	Y=.500*D1/NMAXT	00012700
	CALL TZ(Y,TANS)	00012710
	R=NMAXT*TANS	00012720
	GO TO 240	00012730
230	R=0.0	00012740
240	Y=.7357600*D1/X	00012750
	CALL TZ(Y,TANS)	00012760
	S=1.359100*X*TANS	00012770
	IF(R.GT.S)GO TO 280	00012780
	NU=1+IDINT(S)	00012790
	GO TO 290	00012800
280	NU=1+IDINT(R)	00012810
290	M=0	00012820
	FL=1.	00012830
	LIMIT=(NU/2)	00012840
320	M=M+1	00012850
	FL=FL*(M+A)/(M+1.00)	00012860
	IF(M.LT.LIMIT)GO TO 320	00012870
	M=2*M	00012880
	R=0.0	00012890
	S=0.0	00012900
390	DENOM=2.*(A+N)/X-R	00012910
	IF(DABS(DENOM).LE.1.0D-15)DENOM=DENOM+1.0D-15	00012920
430	R=1./DENOM	00012930
	NMOD2=MOD(N,2)	00012940
	IF(NMOD2.NE.0)GO TO 480	00012950
	FL=FL*(N+2.00)/(N+2.*A)	00012960
	FLMRDA=FL*(N+A)	00012970
	GO TO 490	00012980
480	FLMRDA=0.0	00012990
490	S=R*(FLMRDA+S)	00013000
	IF(N.LE.NMAXT)RR(N)=S	00013010

	N=N-1	00013020
	IF(N.GE.1)GO TO 390	00013030
	FJ(1)=SUM/(1.+S)	00013040
	IF(NMAXT.EQ.0)GO TO 570	00013050
	DO 560 N=1,NMAXT	00013060
560	FJ(N+1)=RR(N)*FJ(N)	00013070
570	DO 640 N=1,NMPL	00013080
	IF(DABS((FJ(N)-FJAPRX(N))/FJ(N)).LE.EPSLON)GO TO 640	00013090
	DO 610 M=1,NMPL	00013100
610	FJAPRX(M)=FJ(M)	00013110
	NIJ=NIJ+5	00013120
	GO TO 290	00013130
640	CONTINUE	00013140
	IF(NMAX.GE.0)RETURN	00013150
	FJ(2)=2.*A*FJ(1)/X-FJ(2)	00013160
	IF(NMAXAB.EQ.1)RETURN	00013170
	DO 650 N=2,NMAXAB	00013180
650	FJ(N+1)=2.*(A-N)*FJ(N)/X-FJ(N-1)	00013190
	RETURN	00013200
	END	00013210

	SUBROUTINE BSCJZ(X,Y,U,V,NMAX,A,ND,IERR,UAPPRX,VAPPRX,RF1,RP2)	BSC13220
	IMPLICIT REAL*8 (A-H,O-Z)	BSC13230
	DIMENSION U(100),V(100),UAPPRX(100),VAPPRX(100),RR1(100),	BSC13240
1	RP2(100)	BSC13250
	IF(A.GE.0.0)GO TO 40	BSC13260
	IERR=1	BSC13270
	RETURN	BSC13280
40	IF(A.LT.1)GO TO 70	BSC13290
	IERR=2	BSC13300
	RETURN	BSC13310
70	IF(X.GT.0.0)GO TO 110	BSC13320
	IF(DABS(Y).LE.1.00-14)GO TO 90	BSC13330
	GO TO 110	BSC13340
90	IERR=3	BSC13350
	RETURN	BSC13360
110	IF(NMAX.GE.0)GO TO 140	BSC13370
	IERR=4	BSC13380
	RETURN	BSC13390
140	IERR=0	BSC13400
	EPSLON=.5D0*10.**(-ND)	BSC13410
	NMP1=NMAX+1	BSC13420
	DO 200 N=1,NMP1	BSC13430
	UAPPRX(N)=0.0	BSC13440
200	VAPPRX(N)=0.0	BSC13450
	Y1=DABS(Y)	BSC13460
	RZ=X**2+Y**2	BSC13470
	RZ=DSQRT(RZ)	BSC13480
	IF(DABS(X).LE.1.00-14)GO TO 290	BSC13490
	PHI=DATAN2(Y1,X)	BSC13500
	IF(X.LT.0.0) PHI=3.141592653589793D0 + PHI	BSC13510
	GO TO 300	BSC13520
290	PHI=1.570796326794896D0	BSC13530
300	C=DEXP(Y1)*(RZ/2.0)**A/DCOS(A*(1.+A))	BSC13540
	SUM2=A*PHI-X	BSC13550
	SUM1=C*DCOS(SUM2)	BSC13560
	SUM2=C*DSIN(SUM2)	BSC13570
	D1=2.3026D0*ND+1.2867D0	BSC13580
	IF(NMAX.GT.0)GO TO 380	BSC13590
	P=0.0	BSC13600
	GO TO 390	BSC13610
380	PARAM=.5D0*D1/NMAX	BSC13620
	CALL TZ(PARAM,TANS)	BSC13630
	R=NMAX*TANS	BSC13640
390	S=1.3591D0*RZ	BSC13650
	PARAM=.73576D0*(D1-Y1)/RZ	BSC13660
	CALL TZ(PARAM,TANS)	BSC13670
	IF(Y1.LT.D1)S=S*TANS	BSC13680
	IF(R.GT.S)GO TO 450	BSC13690
	NJ=1+IDINT(S)	BSC13700
	GO TO 460	BSC13710
450	NJ=1+IDINT(R)	BSC13720
460	N=0	BSC13730
	FL=1.	BSC13740
	C1=1.	BSC13750
	C2=0.	BSC13760
500	N=N+1	BSC13770
	FL=FL*(N+2.*A)/(N+1.00)	BSC13780
	C=-C1	BSC13790
	C1=C2	BSC13800
	C2=C	BSC13810
	IF(N.LT.NJ)GO TO 500	BSC13820

```

R1=0.0
R2=0.0
S1=0.0
S2=0.0
610 C=(2.*(A+N)-X*R1+Y1*P2)**2+(X*P2+Y1*R1)**2
R1=(2.*(A+N)*X-R22*P1)/C
R2=(2.*(A+N)*Y1+R22*P2)/C
FL=FL*(N+1.D0)/(N+2.*A)
C=2.*(N+A)*FL
FLAMB1=C*C1
FLAMB2=C*C2
C=C1
C1=-C2
C2=C
S=R1*(FLAMB1+S1)-P2*(FLAMB2+S2)
S2=R1*(FLAMB2+S2)+P2*(FLAMB1+S1)
S1=S
IF(N.GT.NMAX)GO TO 770
RR1(N)=R1
RR2(N)=R2
770 N=N-1
IF(N.GE.1)GO TO 610
C=(1.+S1)**2+S2**2
U(1)=(SUM1*(1.+S1)+SUM2*S2)/C
V(1)=(SUM2*(1.+S1)-SUM1*S2)/C
IF(NMAX.EQ.0)GO TO 850
DO 840 N=1,NMAX
U(N+1)=RR1(N)*U(N)-RR2(N)*V(N)
840 V(N+1)=RR1(N)*V(N)+RR2(N)*U(N)
850 IF(Y.LT.0.0)GO TO 860
GO TO 880
860 DO 870 N=1,NMP1
870 V(N)=-V(N)
880 DO 950 N=1,NMP1
TEMP1=(U(N)-UAPPRX(N))**2
TEMP1=TEMP1+(V(N)-VAPPRX(N))**2
TEMP1=TEMP1/(U(N)**2+V(N)**2)
IF(TEMP1.LE.EPSLON)GO TO 950
DO 920 M=1,NMP1
JAPPRX(M)=U(M)
920 VAPPRX(M)=V(M)
NII=NI+5
GO TO 460
950 CONTINUE
RETURN
END
SUBROUTINE T2(Y,TANS)
REAL*8 Y,Z,P,TANS,DLOG
IF(Y.GT.10.0)GO TO 40
P=.00005794100*Y-.0017614800
P=Y*P+.020864500
P=Y*P-.12901300
P=Y*P+.8577700
TANS=Y*P+1.1012500
RETURN
40 Z=DLOG(Y)-.77500
P=(.77500-DLOG(Z))/(1.0+Z)
TANS=Y/((1.0+P)*Z)
RETURN
END

```

```

BSC13930
BSC13840
BSC13850
BSC13860
BSC13870
BSC13880
BSC13890
BSC13900
BSC13910
BSC13920
BSC13930
BSC13940
BSC13950
BSC13960
BSC13970
BSC13980
BSC13990
BSC14000
BSC14010
BSC14020
BSC14030
BSC14040
BSC14050
BSC14060
BSC14070
BSC14080
BSC14090
BSC14100
BSC14110
BSC14120
BSC14130
BSC14140
BSC14150
BSC14160
BSC14170
BSC14180
BSC14190
BSC14200
BSC14210
BSC14220
BSC14230
BSC14240
BSC14250
BSC14260
BSC14270
BSC14280
BSC14290
BSC14300
BSC14310
BSC14320
BSC14330
BSC14340
BSC14350
BSC14360
BSC14370
BSC14380
BSC14390
BSC14400
BSC14410
BSC14420

```

Reproduced from
best available copy.

REFERENCES

1. Baum, C.E., "On the Singularity Expansion Method for the Solution of Electromagnetic Interaction Problems," Interaction Note 88, Air Force Weapons Laboratory, Kirtland AFB, NM, December 1971.
2. Taylor, C. D., Crow, T. T., and Chen, K-T, "On the Singularity Expansion Method Applied to Aperture Penetration: Part I Theory," Interaction Note 134, Air Force Weapons Laboratory, Kirtland AFB, NM, May 1973.
3. Rahmat-Samii, Y. and Mittra, R., "Integral Equation Solution and RCS Computation of a Thin Rectangular Plate," Interaction Note 156, Air Force Weapons Laboratory, Kirtland AFB, NM, December 1973.
4. Bladel, J. Van, Electromagnetic Fields, McGraw-Hill, New York, pp. 385-387, 1964.
5. Baum, C. E., "Interaction of Electromagnetic Fields with any Object which has an Electromagnetic Symmetry Plane," Interaction Note 63, Air Force Weapons Laboratory, Kirtland AFB, NM, March 1971.
6. Butler, C. M., "Integral Equation Solution Methods," in Wire Antennas and Scatterers, Short Course Notes, University of Mississippi, April 1972. (See also IEEE Trans., v. AP-20, pp. 731-736, 1972.)
7. Umashankar, K. R., The Calculation of Electromagnetic Transient Currents on Thin Perfectly Conducting Bodies Using the Singularity Expansion Method, Ph.D. Thesis, University of Mississippi, August 1974, pp. 33-34. (See also F. M. Tesche, IEEE Trans., v. AP-21, No. 1, pp. 53-62, 1973.)
8. Davis, W. A., Numerical Solutions to the Problem of Electromagnetic Radiation and Scattering by a Finite Cylinder, Ph.D. Thesis, University of Illinois, 1974.
9. Dunaway, O. C., A Numerical Solution for the Distribution of Time-Harmonic Electromagnetic Fields in an Arbitrary Shaped Aperture in a Ground Screen, M.S. Thesis, University of Mississippi, 1974.